

Light-shining-through-walls via virtual minicharged particles in a magnetic field

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in collaboration with H. Gies, N. Neitz and F. Karbstein



FSU Jena & Helmholtz Institute Jena



arXiv: not yet ☹ ⇒ work in progress! ☺

7th Patras meeting

☀ Mykonos ☀

June 27th, 2011

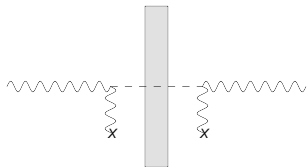
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Light-shining-through-walls scenarios – an overview

- well-known LSW with *real* axions/ALPs:

$$\mathcal{L}_{\text{int,P}} = \frac{1}{4} g_{\text{P}} \phi F_{\mu\nu} \tilde{F}^{\mu\nu}$$



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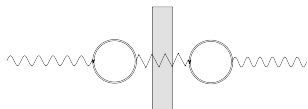
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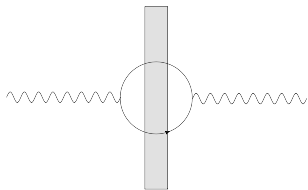
- $\mathcal{L}_{\text{int}} = -\frac{1}{2} \chi F_{\mu\nu} B^{\mu\nu} + e \bar{\psi} \not{A} \psi + e_h \bar{h} \not{B} h$

$$\text{diagonalize} \Rightarrow \mathcal{L}_{\text{effective}} \sim \epsilon e \bar{h} \not{A} h$$

LSW with *real* hidden photons, infer on MCPs indirectly. two couplings: hidden photon e_h & minicharge ϵe



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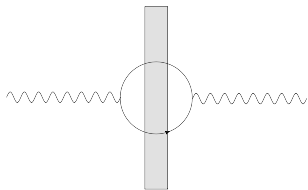
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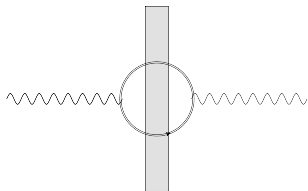
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For zero-field $B = 0$ this was shown to be noncompetitive with current bounds [Jaeckel,Gies 09]

The presented LSW scenarios all have $B \neq 0$
 \Rightarrow How about switching on a magnetic field?
(probably easier in experiment than in theory... ;)

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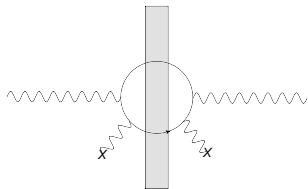
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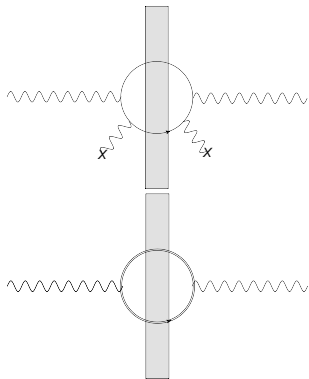
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Prerequisites – going beyond the zero-field limit

object of interest \rightarrow the polarization tensor:
naive ansatz: perturbative treatment of the
in the coupling strength parameter $\frac{2\epsilon e B}{m^2}$
ok in QED but not for MCPs!: both
regimes necessary $\frac{\epsilon e B}{m^2} \ll 1$ **and** $\frac{2\epsilon e B}{m^2} \gg 1$



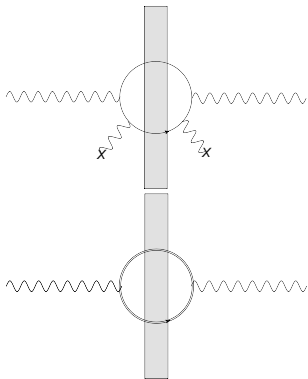
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virtual process: need *full momentum dependence* within polarization tensor

$\theta = \sphericalangle(\vec{B}, \vec{k}) = 0$ & proper time integral
managable! \Rightarrow our choice (only chance)
[see talk by F.Karbstein on thursday!]

Polarization states at zero incidence

States of polarization for $\theta = \angle(\vec{B}, \vec{k}) = 0$

Polarization states: \perp to B , & tilted: recall: $B \neq 0 \Rightarrow v \neq c$
(quasi-longitudinal – in transversal subspace!) $\rightarrow \perp$ [TsEr]
following results are for the \perp mode alone! (work in progress)

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NB: Physics of the \perp propagation mode

number of photons in the \perp mode has yet to be determined:
(what happens for $B = 0$ to $B \neq 0$?) \Rightarrow Our strategy today:
Assume maximum coupling to \perp in the following and report again
with full result!

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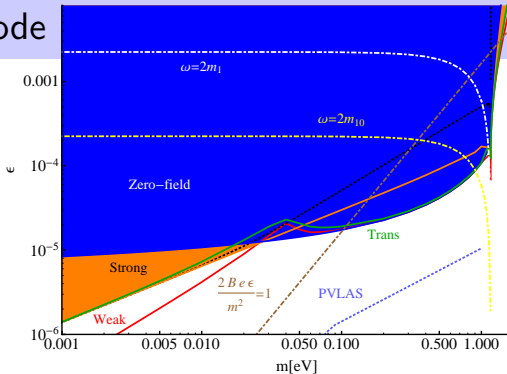
Exclusion bounds for \perp mode

→ $B = 5T, \omega = 2.33\text{eV}$

→ $n_{\text{in}}/n_{\text{out}} = 10^{25}$

(as, e.g. at ALPS)

→ barrier thickness: 1mm, $1\mu\text{m}$



strong B \leftrightarrow weak B

$$\frac{2e\epsilon B}{m^2} \gg 1 \leftrightarrow \frac{2e\epsilon B}{m^2} \ll 1$$

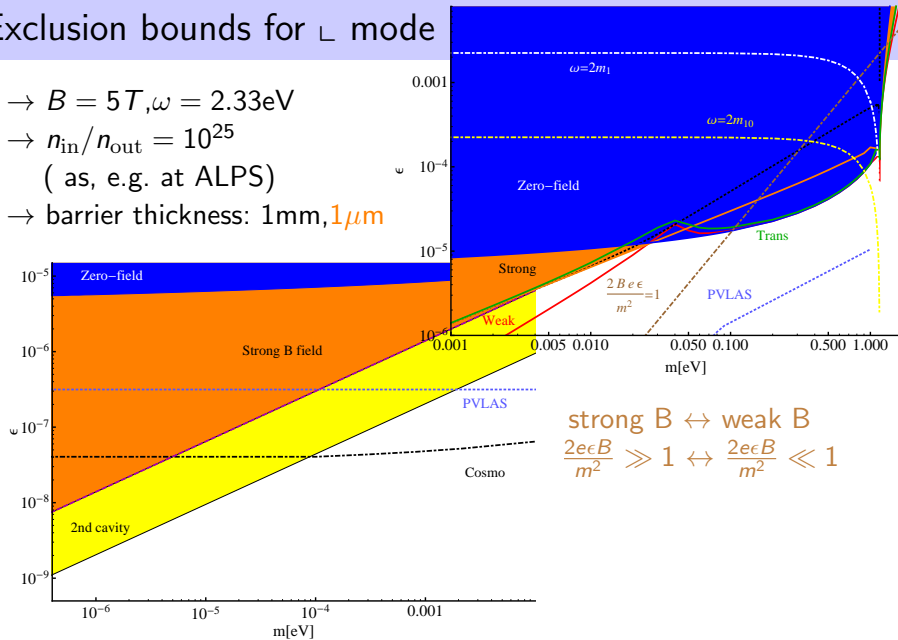
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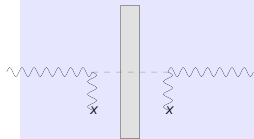
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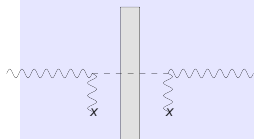
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Physics of virtual LSW – dependence on wall **thickness**

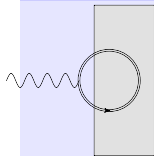


tree-level process: conversion at some point in front of the barrier and reconversion at some point behind the barrier, \Rightarrow **independent** of wall thickness

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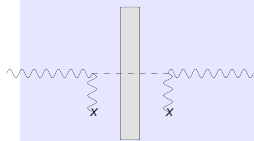


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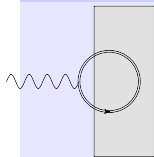


barrier-traversal with **virtual particles:** Compton wavelength smaller than thickness \Rightarrow process obstructed

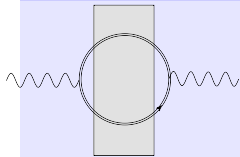
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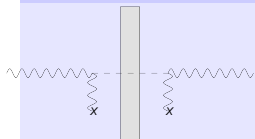


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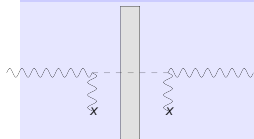
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Physics of **low mass enhancement** (0th Landau level)

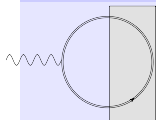


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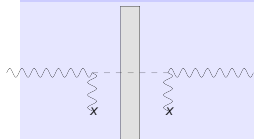


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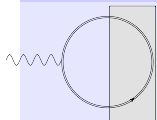


barrier-traversal with **virtual particles:**
Compton wavelength large enough in principle but no barrier traversal

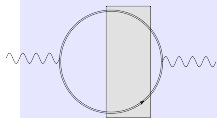
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same "loop position" but larger Compton wavelength. Available "phase space" grows for lower masses
 \Rightarrow *one possible* interpretation of the low mass enhancement

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