

# SOLAR PARAPHOTONS

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- Hidden sectors common in BSM theories
  - dynamical SUSY breaking
  - string compactifications
  - etc.
- Usual non-renormalizable interactions:  
suppressed by the high messenger scale

- Portals: renormalizable interactions
  - Higgs portal
  - neutrino portal
  - **kinetic mixing**
  
- **Kinetic mixing** term happens whenever there is a
  - U(1) gauge field in the hidden sector
    - string Grand Unification
    - gauge-mediated SUSY breaking
    - etc.

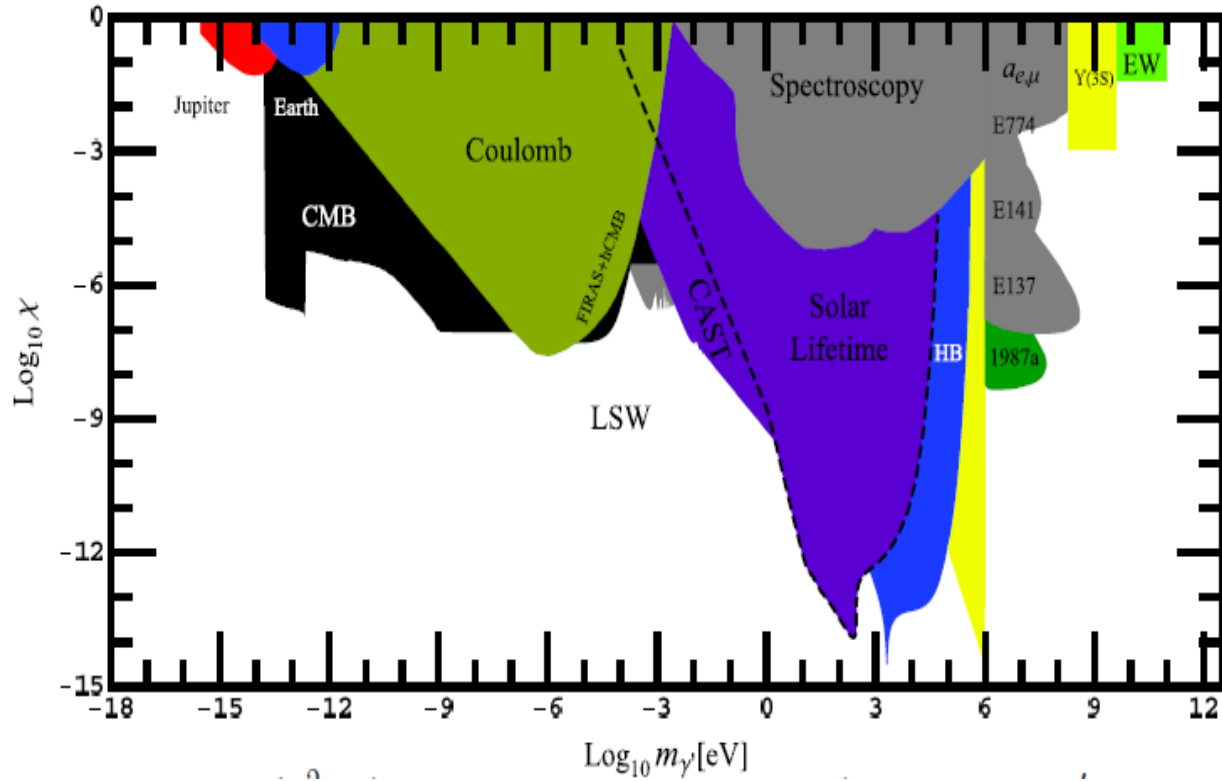
Parameters:

- U(1) hidden gauge coupling
- mixing coupling
- paraphoton mass

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}'^2 - \frac{1}{4}B_{\mu\nu}'^2 - \frac{\sin\chi}{2}F_{\mu\nu}'B_{\mu\nu}' + \frac{m^2}{2}B_{\mu}'^2 - eA_{\mu}'j_{\mu} - gB_{\mu}'\eta_{\mu}.$$

$m, \chi$  relevant for oscillations

constraints on the parameters:



Hidden sector poorly constrained:  
almost every allowed point in  $(m, \chi)$  is motivated

However, three regions of special interest:

1. Mimicking extra neutrino flavour in CMB:

$$(WMAP7: \quad N_{\text{eff}} = 4.34 + 0.86 - 0.88)$$

$$\chi = (1.1 - 2.4) \times 10^{-6}$$

$$m = (10^{-5} - 10^{-2}) \text{ eV}$$

2. TeV scale gravity + strings (some models):

$$M_{\text{Planck}} \text{ related to } \chi$$

$$\chi \sim (10^{-12} - 10^{-10})$$

(lower limit from early LHC searches)

3. “Unified” DM, “secluded” DM, “hidden” Higgs:

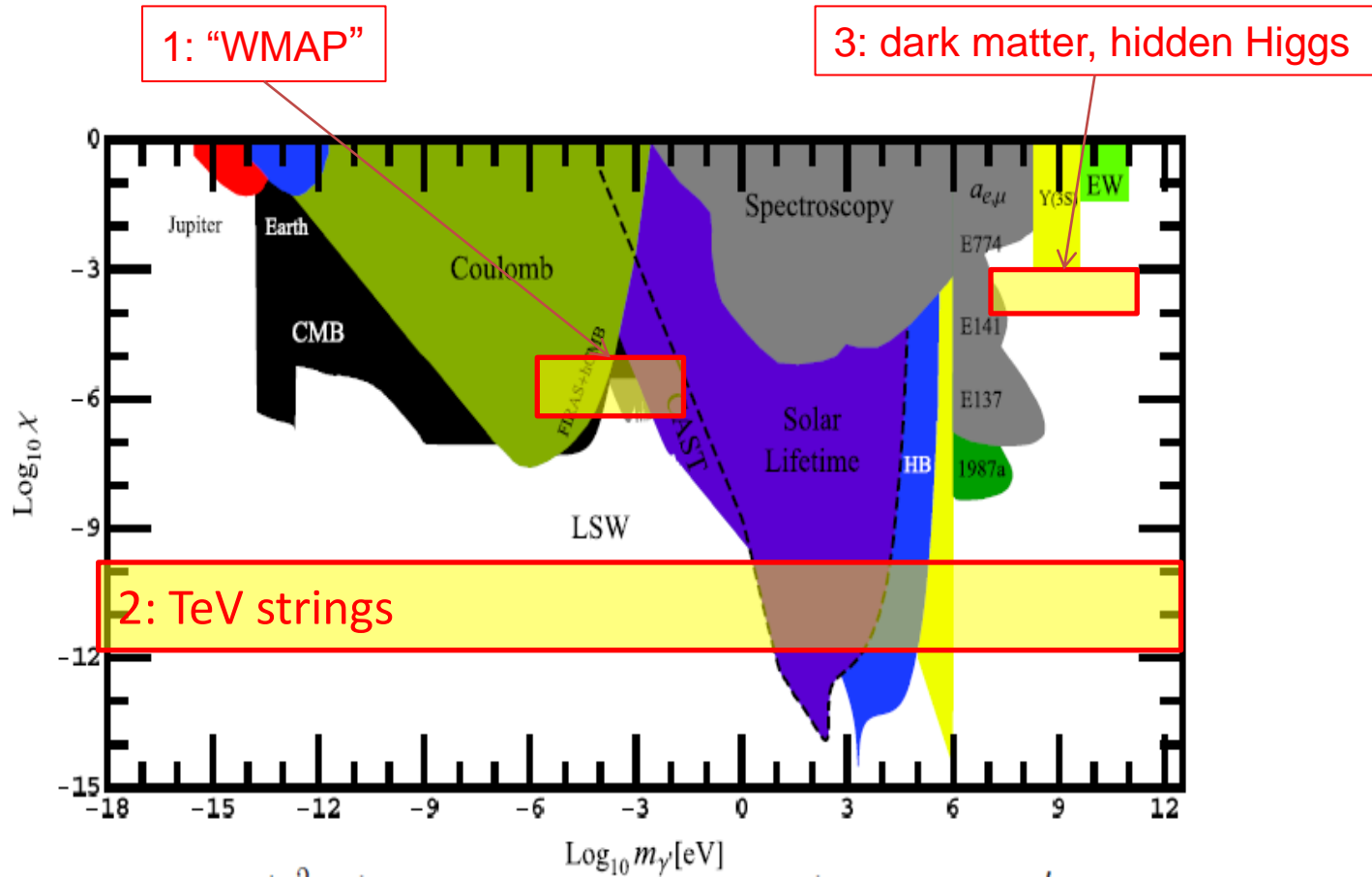
- explain DAMA, PAMELA, ATIC, FERMI, INTEGRAL...

- unusual decays of  $\sim 100$  GeV Higgs (hide it)

$$\chi \sim (10^{-4} - 10^{-3})$$

$$m \sim \text{GeV}$$

# constraints on the parameters:



- 1: LSW, solar eV
- 2: solar keV only
- 3: accelerators

## oscillation probability:

$$P(B) = \left| \sin(\chi - \alpha_f) \cos(\chi - \alpha_i) \exp_{\bar{A}} - \cos(\chi - \alpha_f) \sin(\chi - \alpha_i) \exp_{\bar{B}} \right|^2$$

$$\bar{A}(z) = \bar{A}(0) \exp \left[ -i\omega z + \frac{1}{2\omega} \int_0^z m_{\bar{A}}^2(z') dz' \right] \equiv \bar{A}(0) \exp_{\bar{A}}(z),$$

$$\bar{B}(z) = \bar{B}(0) \exp \left[ -i\omega z + \frac{1}{2\omega} \int_0^z m_{\bar{B}}^2(z') dz' \right] \equiv \bar{B}(0) \exp_{\bar{B}}(z),$$

where:

$$m_{\bar{A}}^2 = \frac{\Pi \cos^2 \alpha - G \sin^2 \alpha}{\cos 2\alpha} = \Pi(1 - \tan \chi \tan \alpha),$$

$$m_{\bar{B}}^2 = \frac{G \cos^2 \alpha - \Pi \sin^2 \alpha}{\cos 2\alpha} = \Pi(1 + \tan \chi \cot \alpha),$$

$$\tan 2\alpha = \frac{\Pi \sin 2\chi}{\Pi \cos 2\chi - m^2}$$

$$\Pi = \Pi(z) = \omega_p^2(z) + i\omega\Gamma(z)$$



observation of astro paraxotons  
from an optically thick emission region:

$$\alpha_f = 0$$

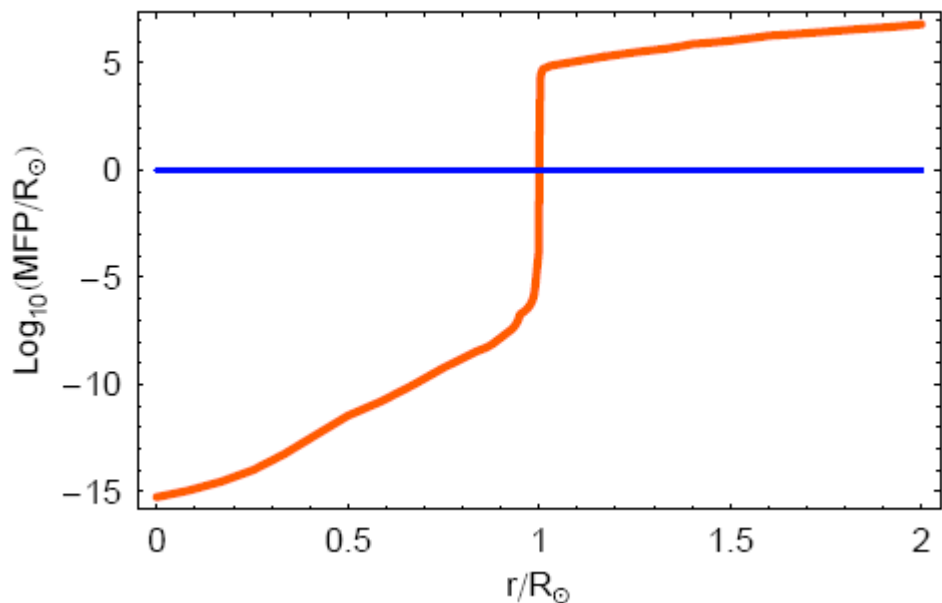
$$\tau \equiv \int \Gamma(z') dz' \gg 1$$

$$P(B) \simeq \cos^2 \chi |\sin(\chi - \alpha_i)|^2 \simeq \chi^2 \frac{m_1^4}{(\omega_{p,i}^2 - m_1^2)^2 + \omega^2 \Gamma_i^2}$$

# SOLAR PARAPHOTONS: converted from thermal photons inside the Sun due to kinetic mixing

Approximation: Sun = optically thick plus optically thin parts

Plot: mean free path of 1 keV photon vs. distance from the solar center



Optically thick:  $0 \leq r/R_{\text{sun}} \leq 0.993$

Transparent:  $0.993 < r/R_{\text{sun}} < 215$  (mostly corona)

Transition region too thin to contribute significantly even at resonance

## SOLAR PARAPHOTONS:

converted from thermal photons inside the Sun due to kinetic mixing

- opaque source: conversion probability determined by the emission point

$$\simeq \chi^2 \frac{m_1^4}{(\omega_{p,i}^2 - m_1^2)^2 + \omega^2 \Gamma_i^2}$$

The diagram shows the conversion probability formula with red arrows pointing from labels in red boxes to the corresponding parts of the equation:

- mixing** points to the  $\chi^2$  term.
- plasma frequency** points to the  $\omega_{p,i}^2$  term in the denominator.
- paraphoton mass** points to the  $m_1^2$  term in the denominator.
- energy** points to the  $\omega^2$  term in the denominator.
- photon absorption** points to the  $\Gamma_i^2$  term in the denominator.

- convolve with the photon emission rate and integrate over the Sun
- contribution of the transparent part subleading for the parameters of interest

## SOLAR PARAPHOTONS:

converted from thermal photons inside the Sun due to kinetic mixing

- convolve with the photon emission rate and integrate over the Sun:

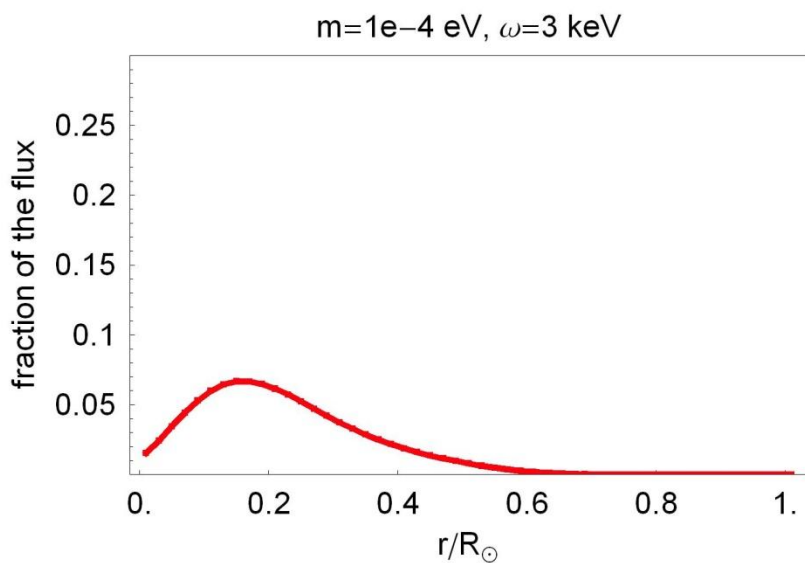
$$\frac{d\Phi}{d\omega} = \frac{3 \times 10^{24}}{\text{cm}^2 \cdot \text{s} \cdot \text{eV}} \left( \frac{\chi}{10^{-5}} \right)^2 \left( \frac{m_1}{\text{eV}} \right)^4 f_1(\omega, m_1)$$

$$f_1(\omega, m_1) = 1 \text{ eV} \times \omega^2 \int_0^1 d\xi \xi^2 \frac{\Gamma(\xi R_\odot)}{e^{\omega/T(\xi R_\odot)} - 1} \frac{1}{(\omega_p^2(\xi R_\odot) - m_1^2)^2 + \omega^2 \Gamma(\xi R_\odot)^2}$$

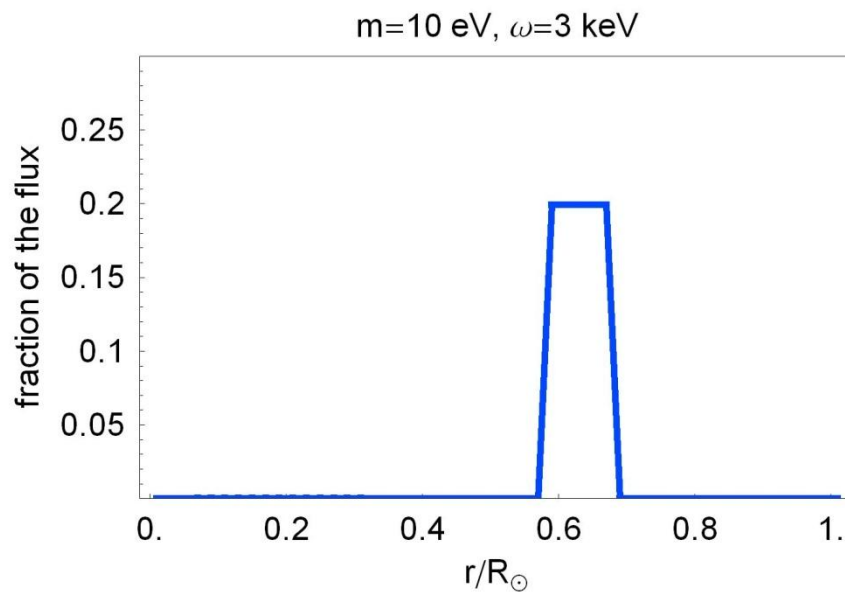
plasma frequency in the Sun 0.1-300 eV:  
if mass in this range, resonance dominates

if no resonance, central keV temperature dominates

# SOLAR PARAPHOTONS: converted from thermal photons inside the Sun due to kinetic mixing

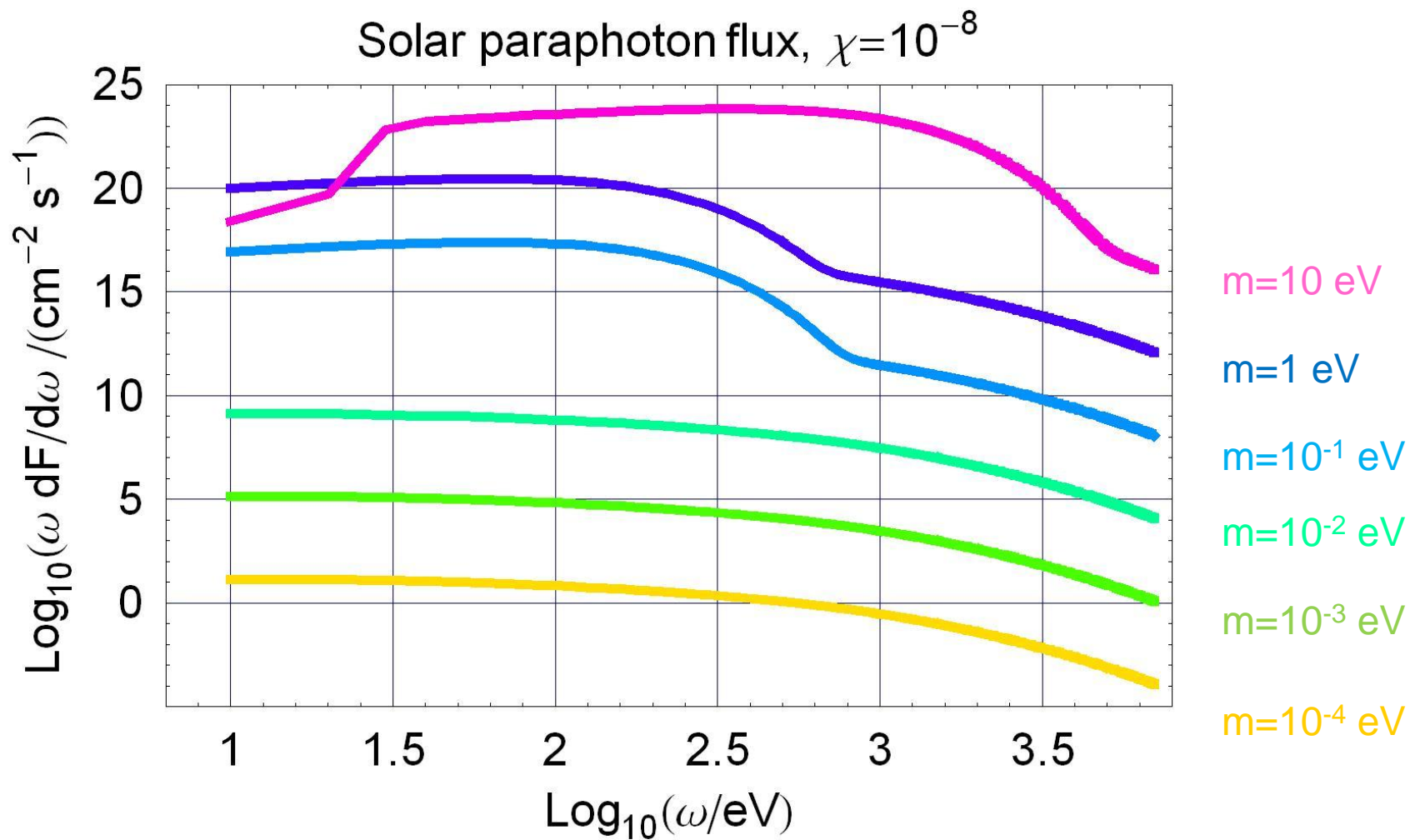


no resonance:  
central part brighter

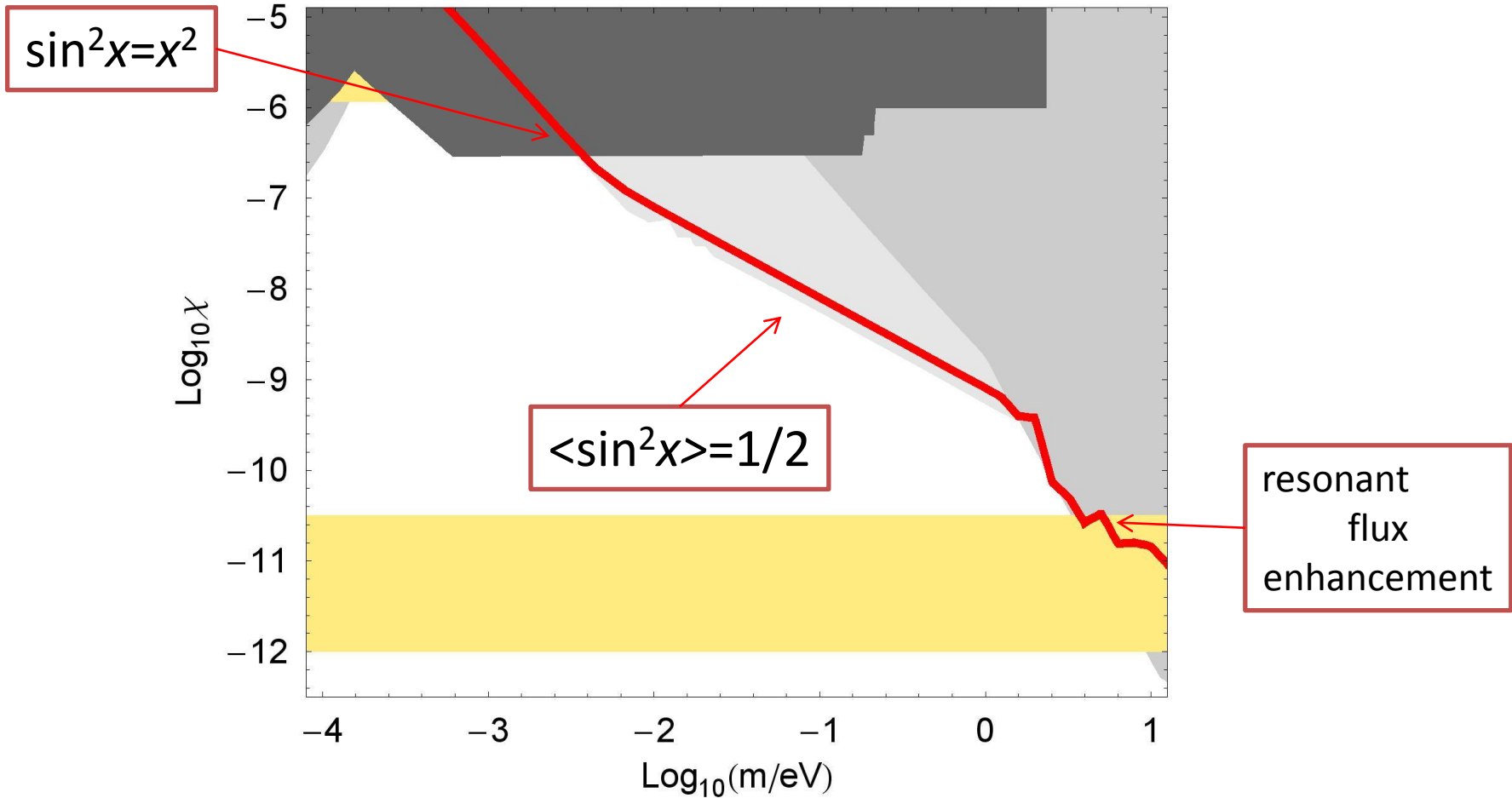


resonance:  
thin slice brighter

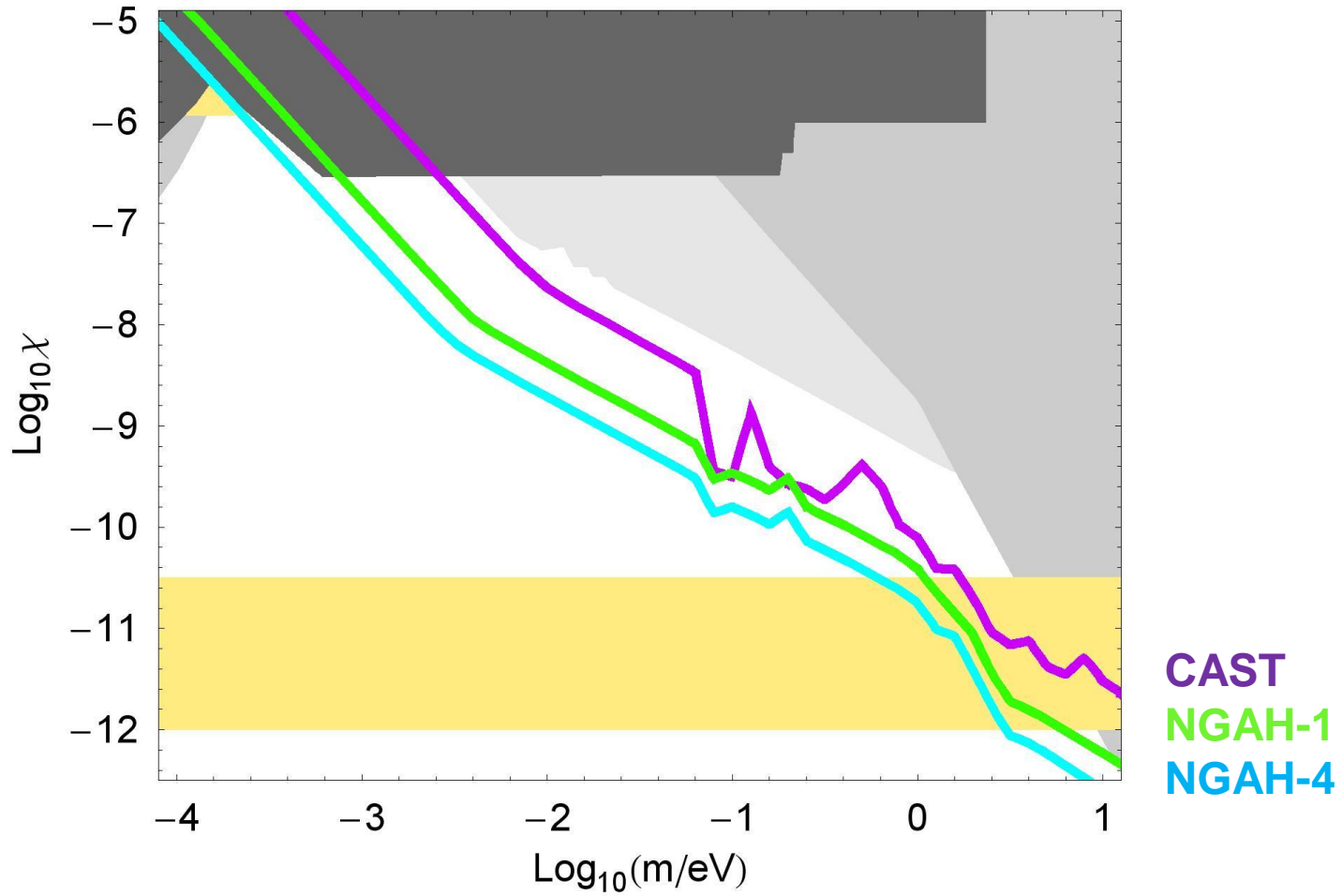
# SOLAR PARAPHOTONS: converted from thermal photons inside the Sun due to kinetic mixing



**DETECTION (vacuum):**  $P_{\text{vac}} = 4\chi^2 \sin^2 \left( \frac{Lm^2}{2\omega} \right)$



# DETECTION PROSPECTS:





Transparent part:

an approximate relation between  
photon and paraphoton fluxes

$$\frac{d\Phi}{d\omega} = \frac{d\Phi_{\text{obs}}}{d\omega} \frac{\int_{R_{\odot}}^D r^2 dr \frac{n_e^2(r)}{\exp(\omega/T(r))-1} P(B, r)}{\int_{R_{\odot}}^D r^2 dr \frac{n_e^2(r)}{\exp(\omega/T(r))-1} (1 - P(B, r))}$$

where

$$P(B) = \frac{1}{2} - \frac{1}{2} \cos 2\chi \frac{\cos 2\chi |\omega_{p,i}^2 - m_1^2| - \sin^2 2\chi \omega_{p,i}^2}{\sqrt{(\omega_{p,i}^2 - m_1^2)^2 + \omega_{p,i}^4 \tan^2 2\chi}}$$

for keV energies and interesting  $m$ ,  $\chi$   
the corona contribution is negligible

## FLARES:

[assume a homogeneous blob of plasma with thermal emission]

$$\frac{d\Phi}{d\omega} = \frac{d\Phi_{\text{obs}}}{d\omega} \frac{\int_{R_{\odot}}^D r^2 dr \frac{n_e^2(r)}{\exp(\omega/T(r))-1} P(B, r)}{\int_{R_{\odot}}^D r^2 dr \frac{n_e^2(r)}{\exp(\omega/T(r))-1} (1 - P(B, r))}$$

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# FLARES:

[assume a homogeneous blob of plasma with thermal emission]

[no resonance:  $\chi^4$ ]

$$\chi \gtrsim 6 \times 10^{-4} \left( \frac{F_{\text{obs}}}{10^5 \text{ cm}^{-2} \text{ s}^{-1} \text{ eV}^{-1}} \right)^{-1/4} \left( \frac{t}{10^3 \text{ s}} \right)^{-1/8} \left( \frac{n}{10^{-3} \text{ Hz}} \right)^{1/8} \left( \frac{S}{10 \text{ cm}^2} \right)^{-1/4} \left( \frac{\omega}{\text{keV}} \right)^{-1/4}$$

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[resonance:  $\chi^2$ ]

$$\chi \gtrsim 8 \times 10^{-7} \left( \frac{F_{\text{obs}}}{10^5 \text{ cm}^{-2} \text{ s}^{-1} \text{ eV}^{-1}} \right)^{-1/2} \left( \frac{t}{1 \text{ s}} \right)^{-1/4} \left( \frac{n}{10^{-3} \text{ Hz}} \right)^{1/4} \left( \frac{S}{10 \text{ cm}^2} \right)^{-1/2} \left( \frac{\omega}{\text{keV}} \right)^{-1/2}$$

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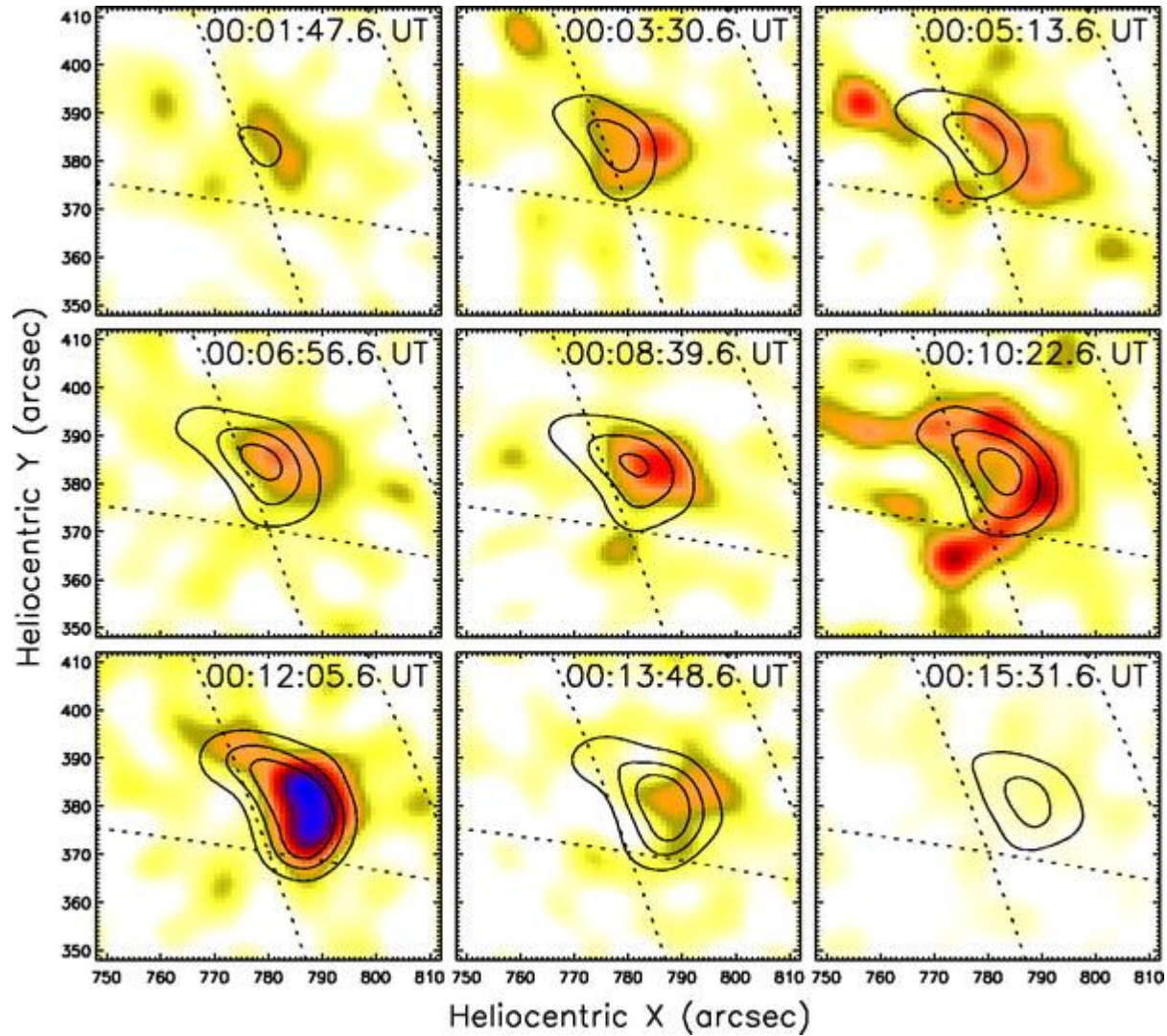
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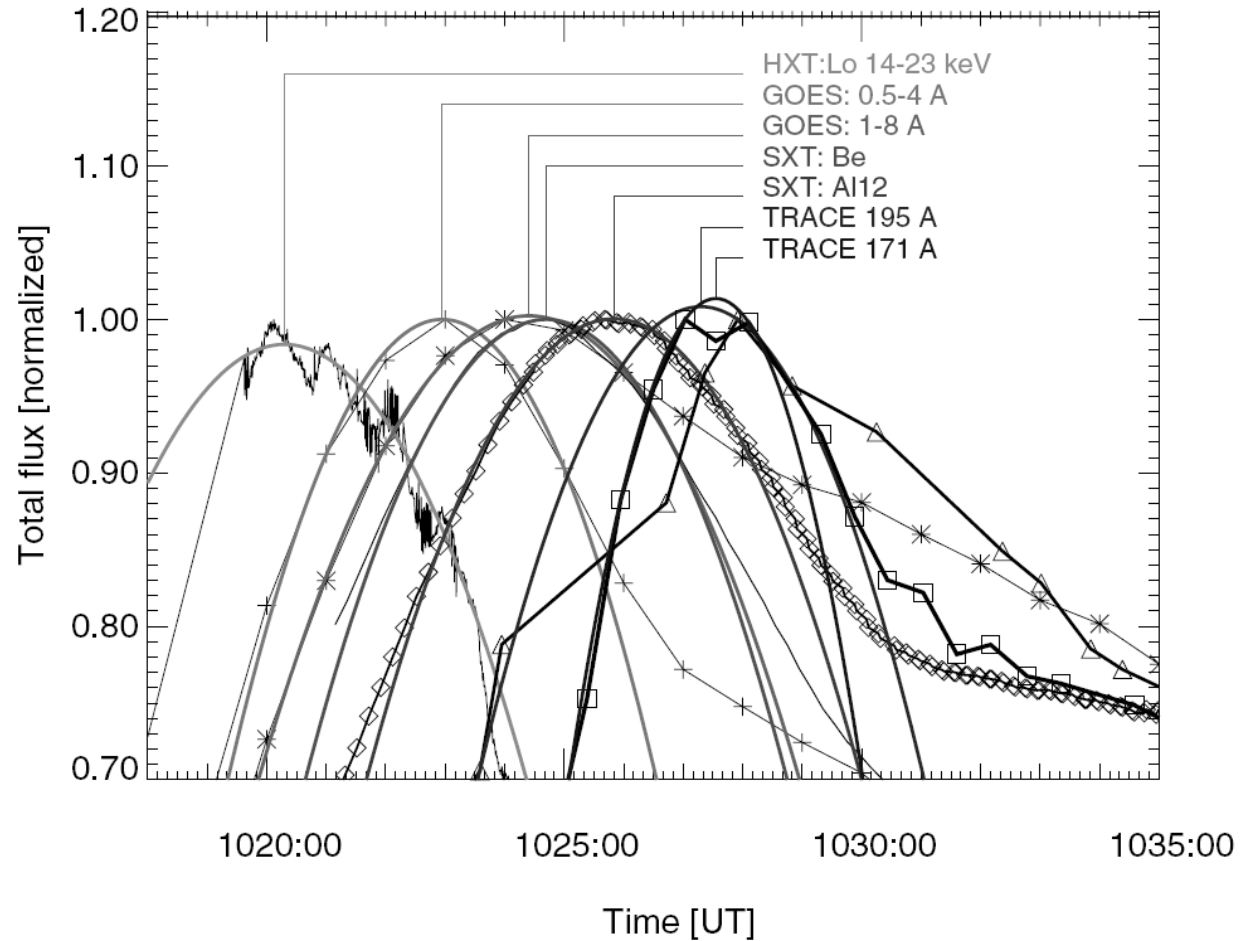
# FLARES: look rapidly evolving

April 14, 2002 event



# FLARES: $n_e$ is NOT rapidly evolving

[reason of variability: temperature change, not movement]



July 14, 2000 event  
[Aschwanden]

# FLARES: $n_e$ is NOT rapidly evolving

April 14, 2002 event

Time	$n = \sqrt{EM/V}$ ( $10^{10} \text{cm}^{-3}$ )
23:57:48-00:04:00	9.3
00:05:00-00:10:00	11.8
00:16:00-00:24:00	8.8

[ApJ 730 (2011) L22]

# FLARES:

[assume a homogeneous blob of plasma with thermal emission]

[resonance:  $\chi^2$ ]

$$\chi \gtrsim 8 \times 10^{-7} \left( \frac{F_{\text{obs}}}{10^5 \text{ cm}^{-2} \text{ s}^{-1} \text{ eV}^{-1}} \right)^{-1/2} \left( \frac{t}{1 \text{ s}} \right)^{-1/4} \left( \frac{n}{10^{-3} \text{ Hz}} \right)^{1/4} \left( \frac{S}{10 \text{ cm}^2} \right)^{-1/2} \left( \frac{\omega}{\text{keV}} \right)^{-1/2}$$

**DISCOVERY, NOT EXCLUSION?**

# FLARES:

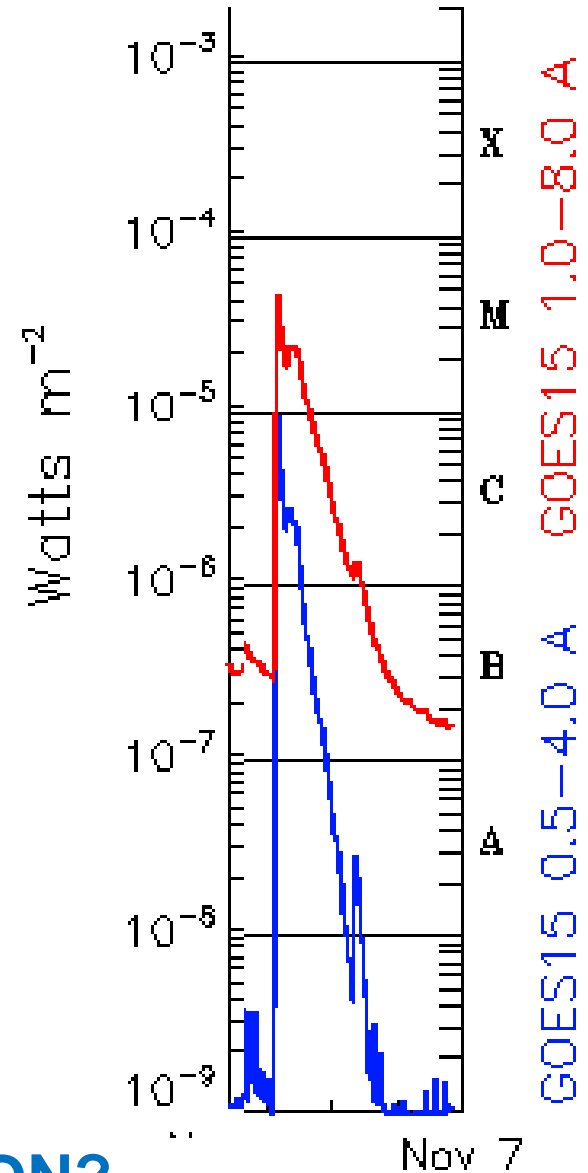
RESONANCE

DROP IN THE REGULAR  
(NORMAL PHOTONS) LIGHT CURVE!

GOOD TIME-RESOLUTION  
SOLAR INSTRUMENTS:

- SOXS (India)
- SPHINX (Russia-Poland)
- ...

**COINCIDENCE =  
FIRM DISCOVERY, NOT EXCLUSION?**



## **CONCLUSIONS:**

1. PARAPHOTONS=WELL-MOTIVATED SM EXTENSION
2. PARASITICAL SEARCH WITH HELIOSCOPIES
3. RESONANCES IN FLARES: A DISCOVERY CHANNEL  
(HELIOSCOPIES + NORMAL TELESCOPES)