

# White Dwarfs as Physics Laboratories: The Axion case

Jordi Isern  
ICE-CSIC/IEEC

Collaboration:

L. Althaus, S. Catalán, A. Córscico, E. García-Berro, M. Salaris, S. Torres

7<sup>th</sup> Patras Workshop on axions, WIMPs & WISPs  
Mykonos, June 26<sup>th</sup>, 2011

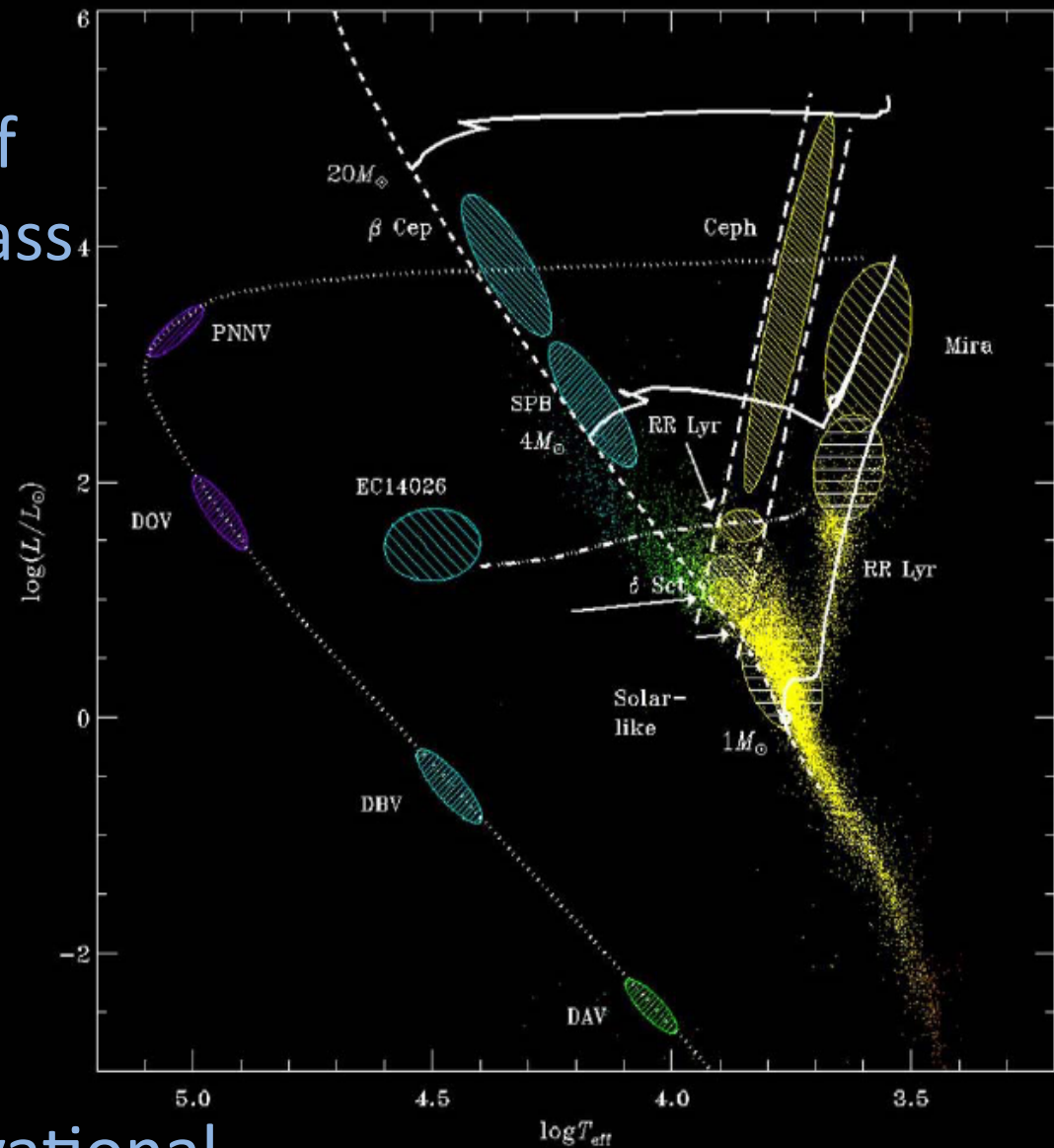
# The white dwarf population is one of the best studied!

# They are the end stage of low and intermediate-mass stars

# Their evolution is just a cooling process

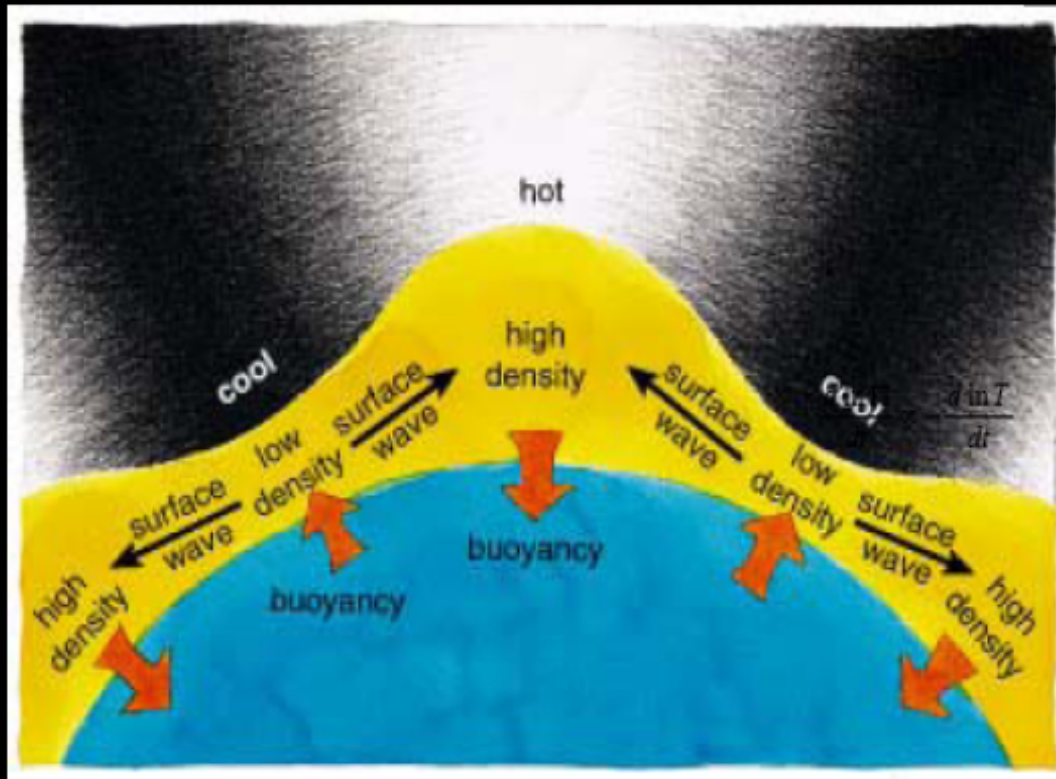
# The basic physical ingredients of their evolution are well identified (not all has been satisfactorily solved yet)

# Impressively solid observational background for testing theory.



Courtesy of Christensen-Dalgaard

# Non-radial g-modes



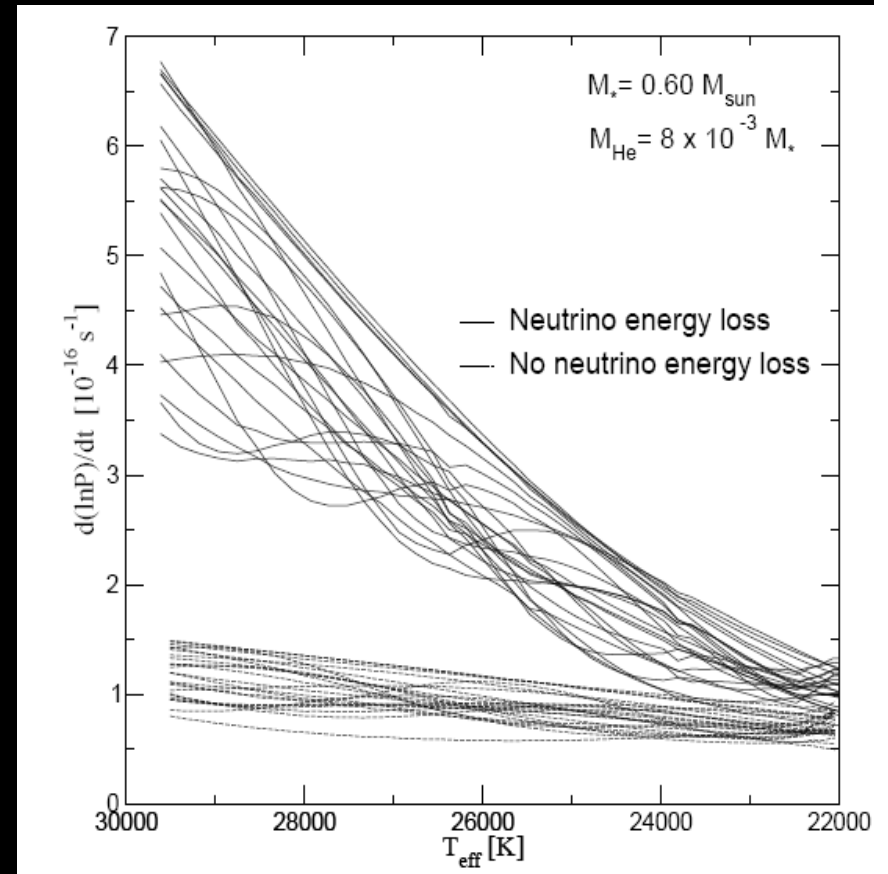
- Long period waves  $\sim 10^2 - 10^3$  s
- Gravity is the restoring force

$$\frac{\dot{P}}{P} = -a \frac{\dot{T}}{T} + b \frac{\dot{R}}{R}$$

# The period increases as the star cools down and decreases as it contracts.

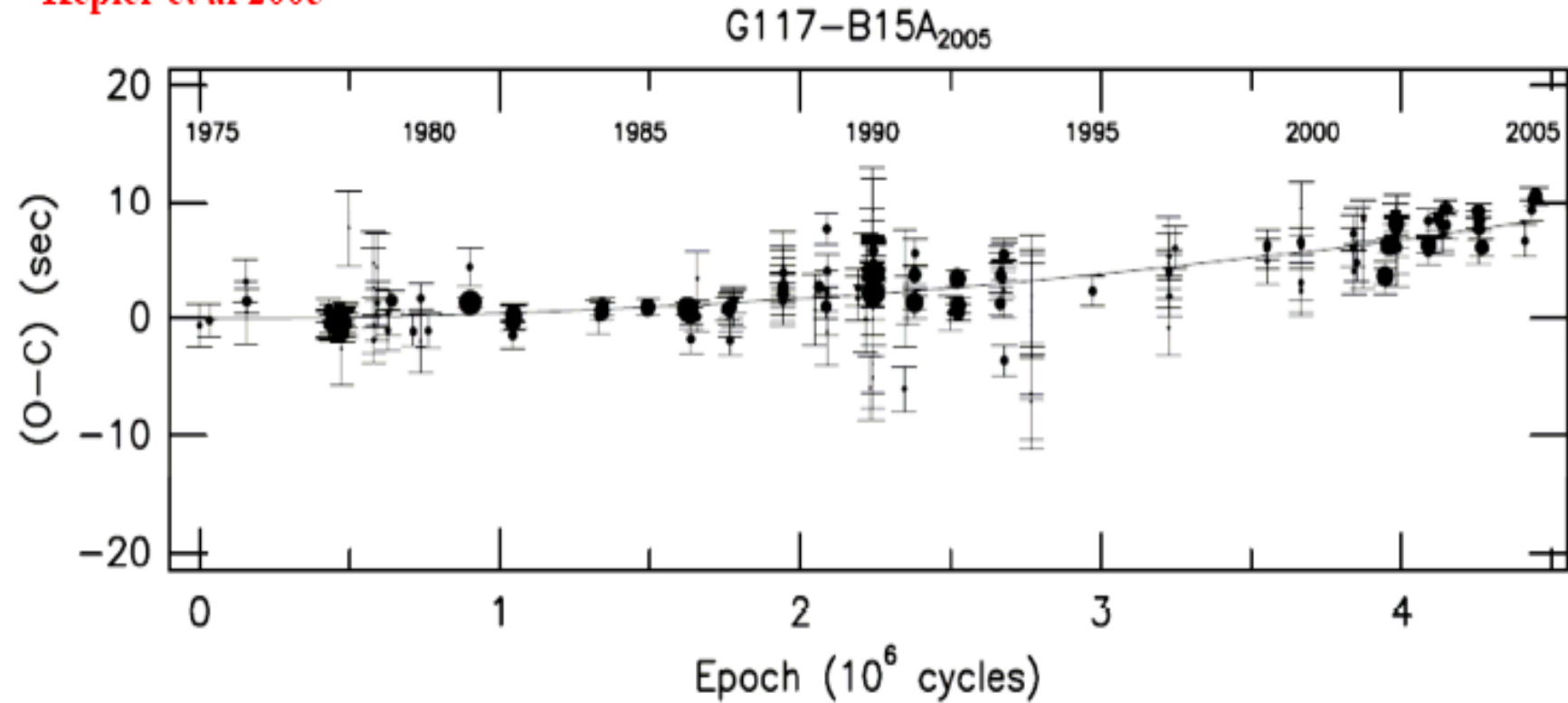
# The radial term can be neglected for cool enough stars (DAV, DBV)

- DOV variables: the drift can be positive or negative depending on the mode
  - PG1159-35:  $P = 516$  s and  $dP/dt = 13.07 \pm 0.3 \times 10^{-11}$  s/s
- DBV variables: the drift is always positive.  $dP/dt \sim 10^{-13} - 10^{-14}$  s/s. No drift measurements
- DAV variables: the drift is always positive.
  - G117-B15A:  $P = 215.2$  s,  $dP/dt = 3.57 \times 10^{-15}$  s/s (Kepler et al 2005)
  - R548:  $P = 213.13$  s,  $dP/dt \leq 5.5 \times 10^{-15}$  s/s



Còrsico and Athaus, 2004

Kepler et al 2005



$$\dot{\Pi} = (12.0 \pm 3.5) \times 10^{-15} \text{ s/s}$$

The first value (Kepler et al'91) was a factor of 2 larger than expected.

Three solutions:

- Observational error
- White dwarfs with "IME" cores
- Exotic source of cooling

$$M_{bol}(t) = -2.5 \log L(t) + ctn$$

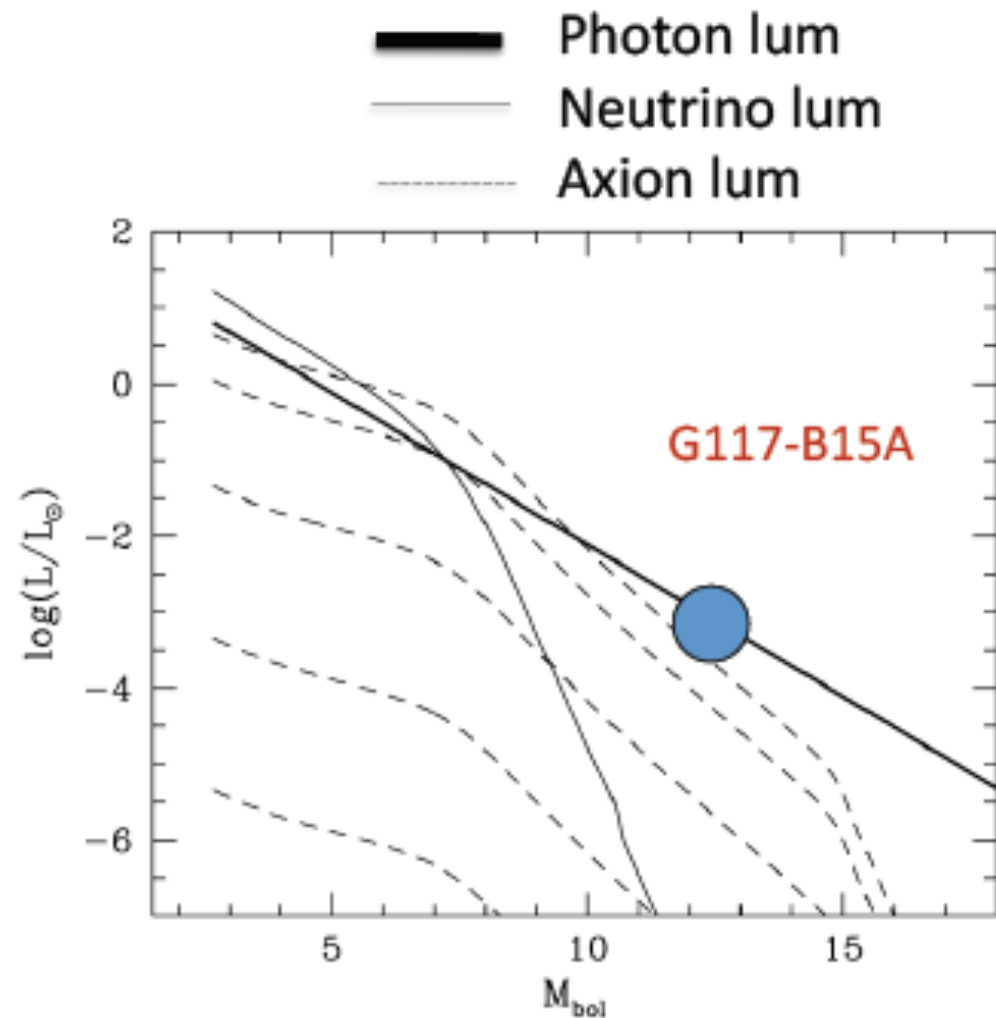
$$\varepsilon_a = 1.08 \cdot 10^{23} \alpha \frac{Z^2}{A} T_7^4 F(\Gamma)$$

$$\alpha = \frac{g_{ae}^2}{4\pi}$$

DFSZ axions

Bremsstrahlung is dominant

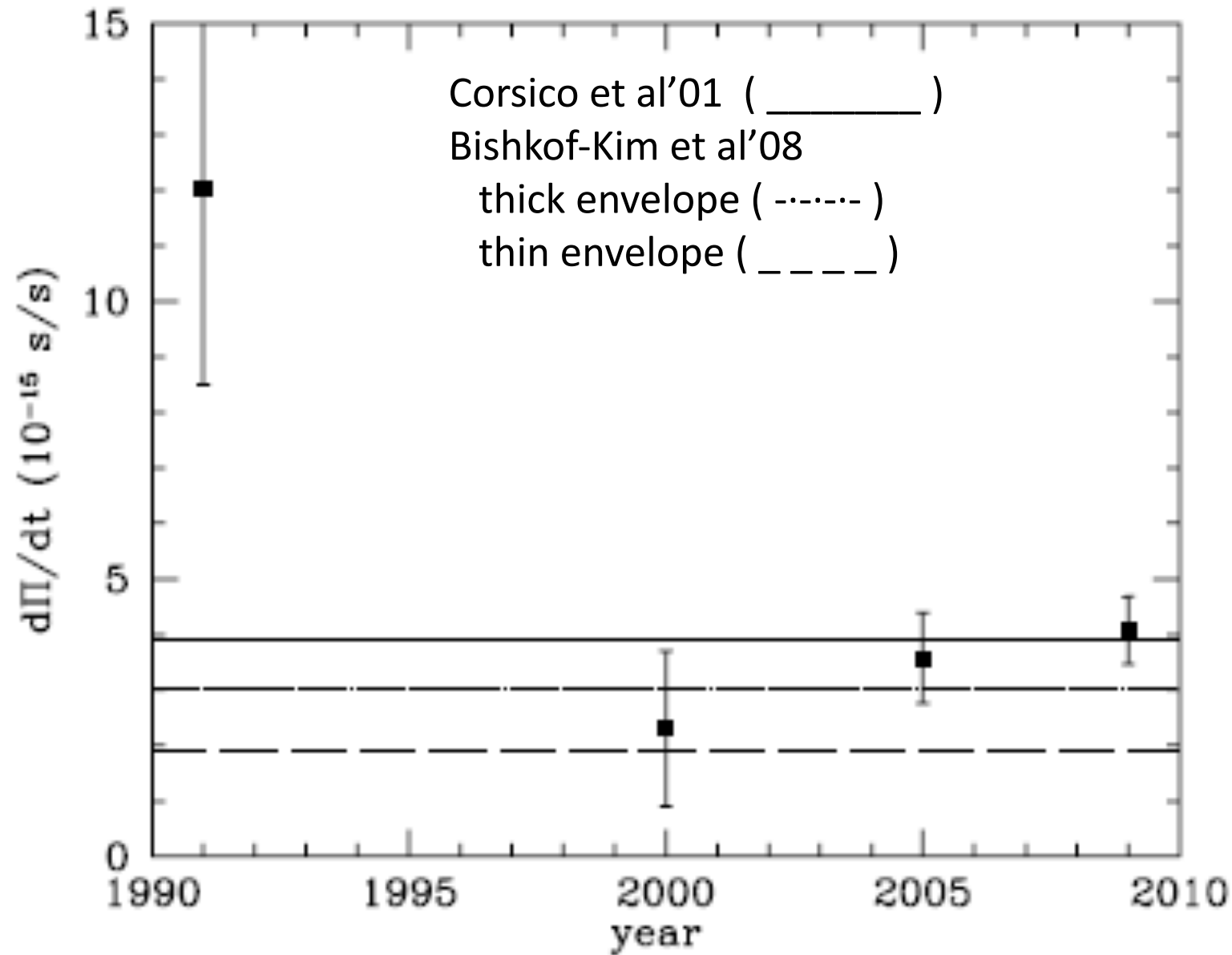
Nakagawa et al 1987, 1988



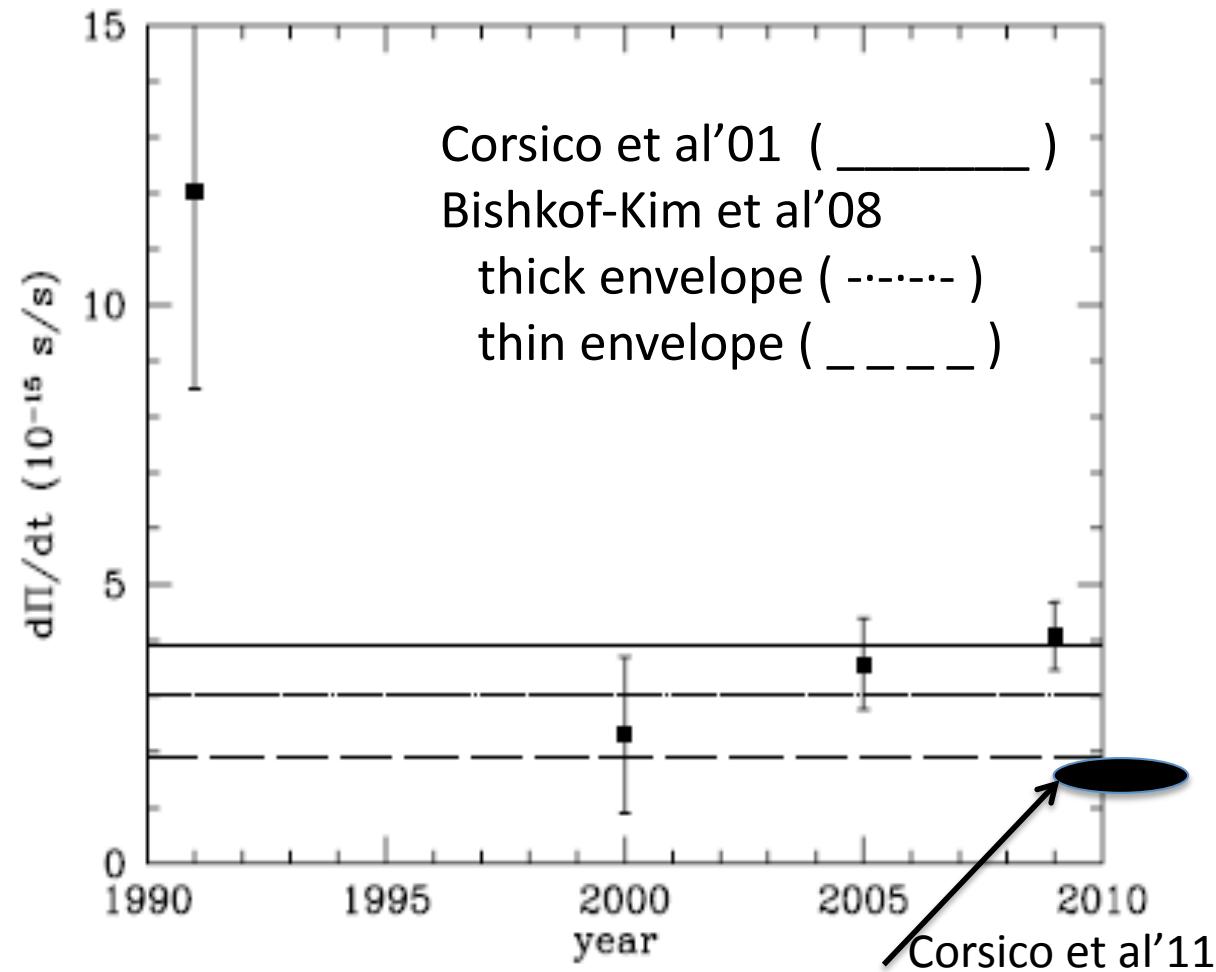
$$g_{ae} \sim 2.2 \times 10^{-13} \quad (m_a \sim 8 \text{ meV}) \quad \text{Isern+'92}$$

# Evolution of the measurements of the period of pulsation period drift of G117-B15A

Isern+'10

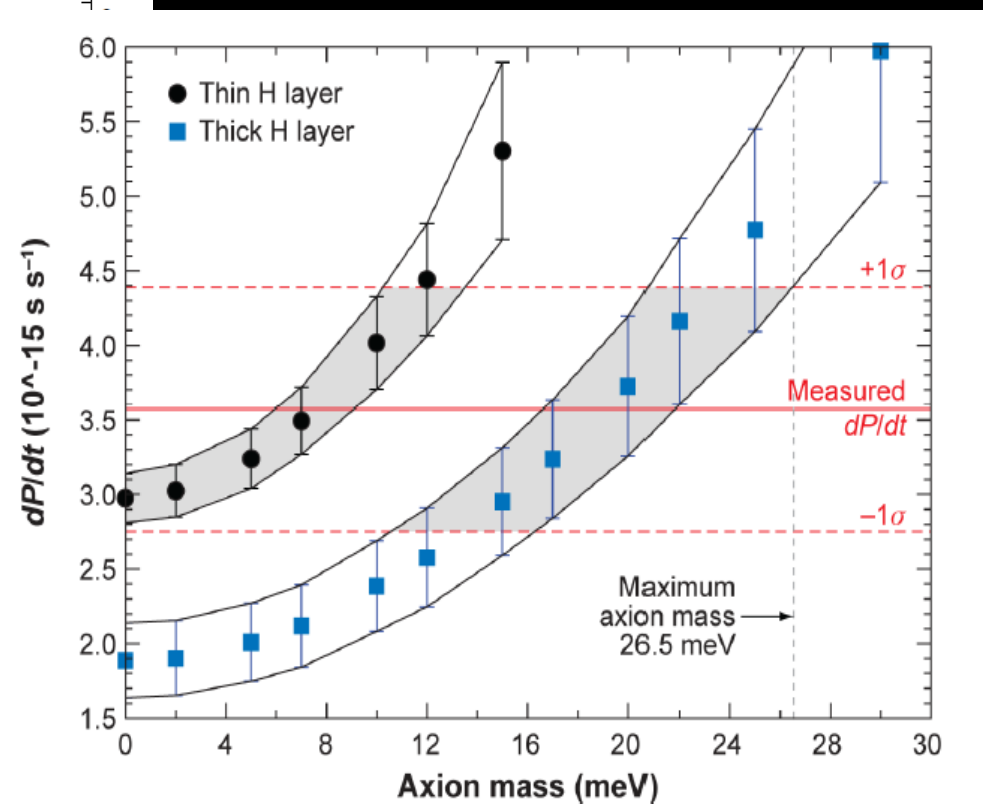
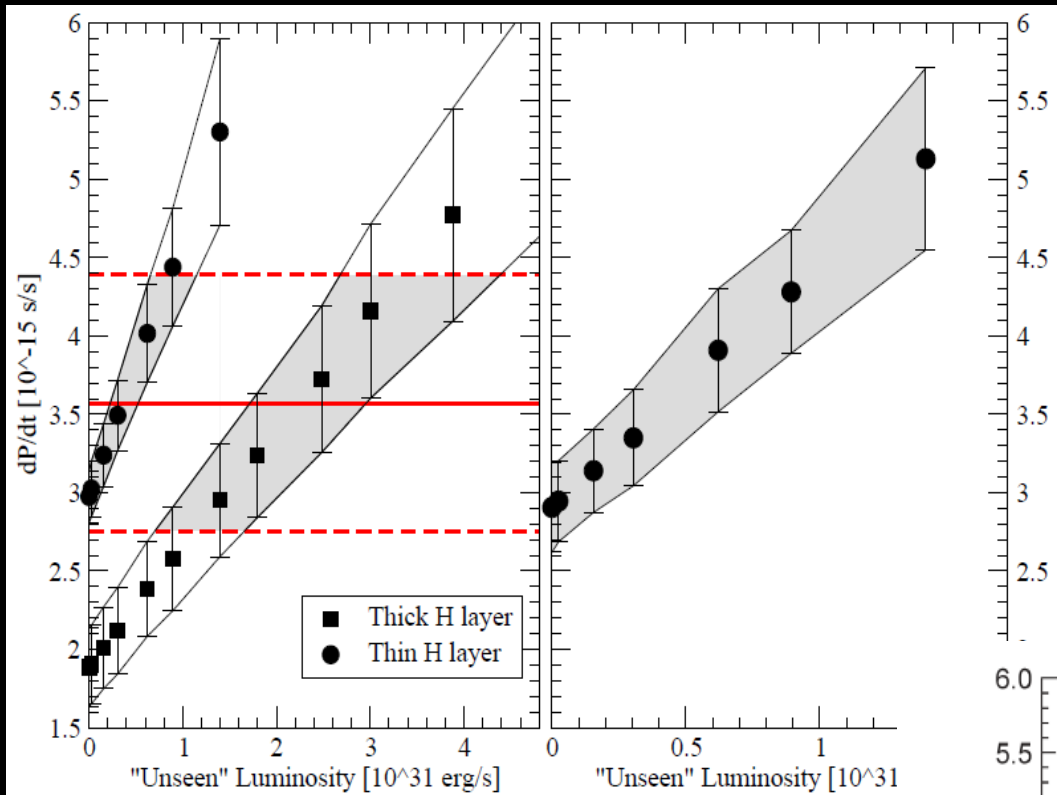


## Observed and predicted secular drift of G117-B15A

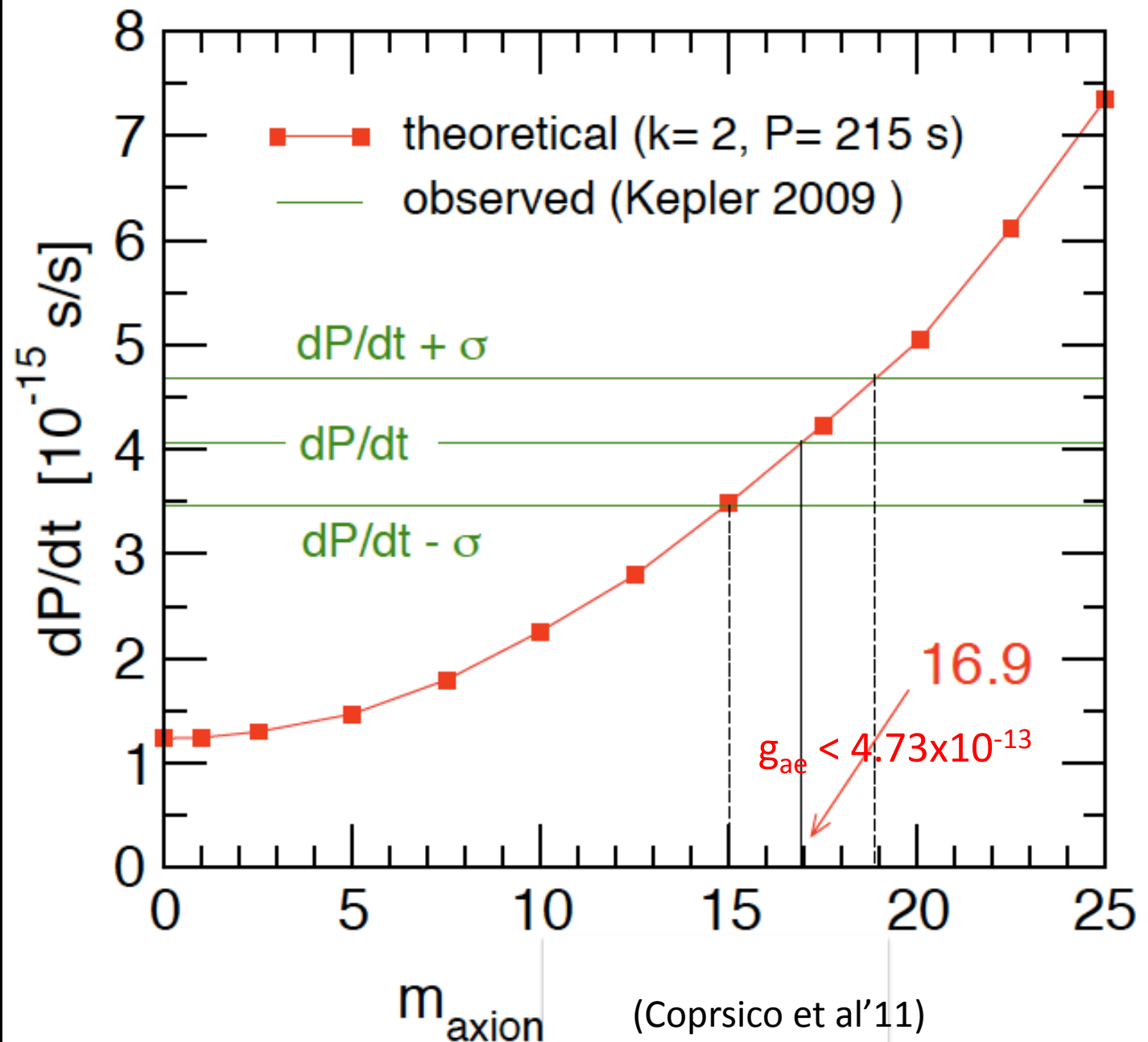


**Comparison of the theoretical predictions (lines) with the observations. An extra cooling term seems necessary.**





Bischoff-Kim et al 2008:  
 $m_a < 13 - 26$  meV  
 $g_{ae} < 3.64 \times 10^{-13} - 7.28 \times 10^{-13}$



# White dwarf cooling

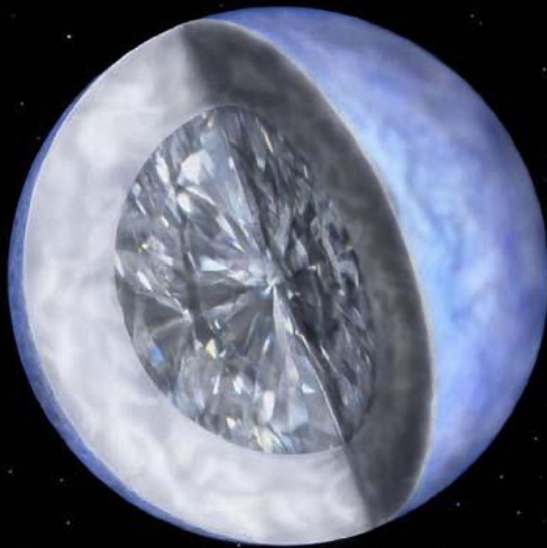
$$L + L_v + (L_e) = - \int_{M_{WD}} c_v \frac{dT_c}{dt} dm - \int_{M_{WD}} T \left( \frac{\partial P}{\partial T} \right)_{V,x} \frac{dV}{dt} dm + (l_s + e_s) \dot{m}_e + (\epsilon_e)$$

A  $L(T_c)$  relationship is necessary to solve this equation

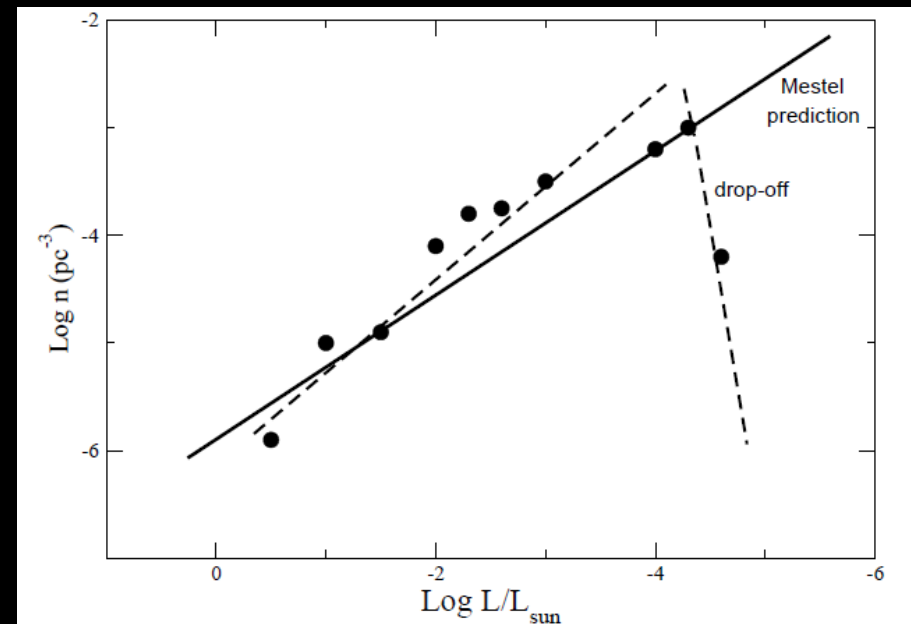
It depends on the properties of the envelope.

$$L \propto T^\alpha$$

$$\alpha \approx 2.5 - 2.7$$



CO.core/He-environment/H-environment



## The luminosity function

Number of white dwarfs per unit of volume and magnitude versus luminosity

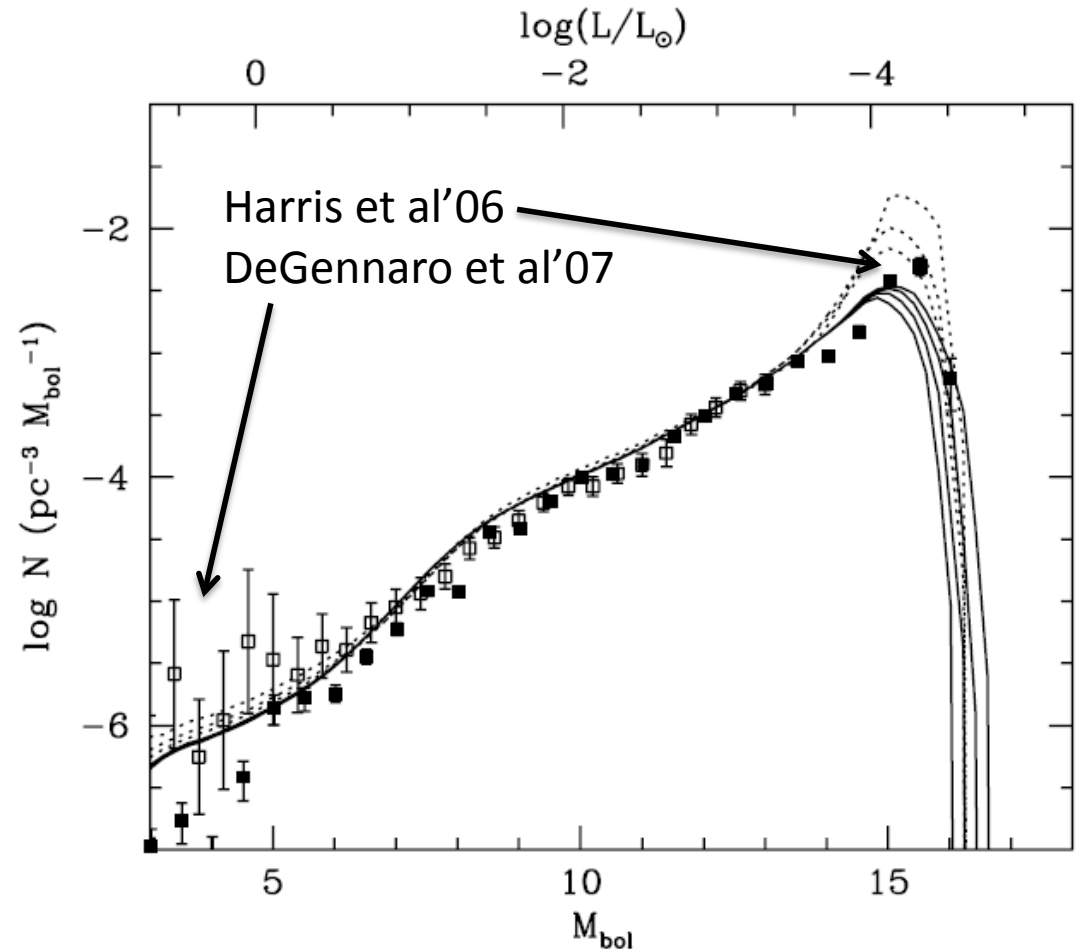
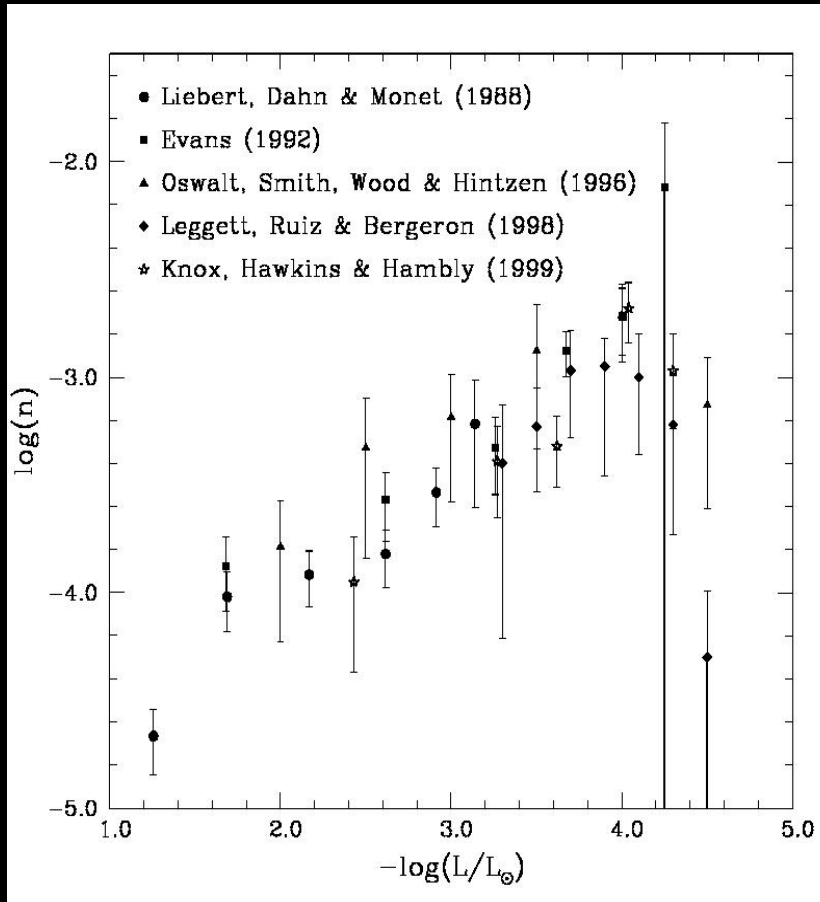
$$n(L) = \int_{M_l}^{M_u} \Phi(M) \Psi(T_G - t_{cool} - t_{ps}) \tau_{cool} dM$$

- 1.-  $n(L)$  is the observed distribution
- 2.-  $\Phi, \Psi$  are the IMF and SFR respectively.  $T_G$  is the age of the Galaxy
- 3.-  $t_{cool}$  is the cooling time  
 $t_{ps}$  is the lifetime of the progenitor  
 $\tau_{cool}$  is the characteristic cooling time  
Hidden an IMFR

If the 3 ingredients are known it is possible to use the WDLF to test new physics

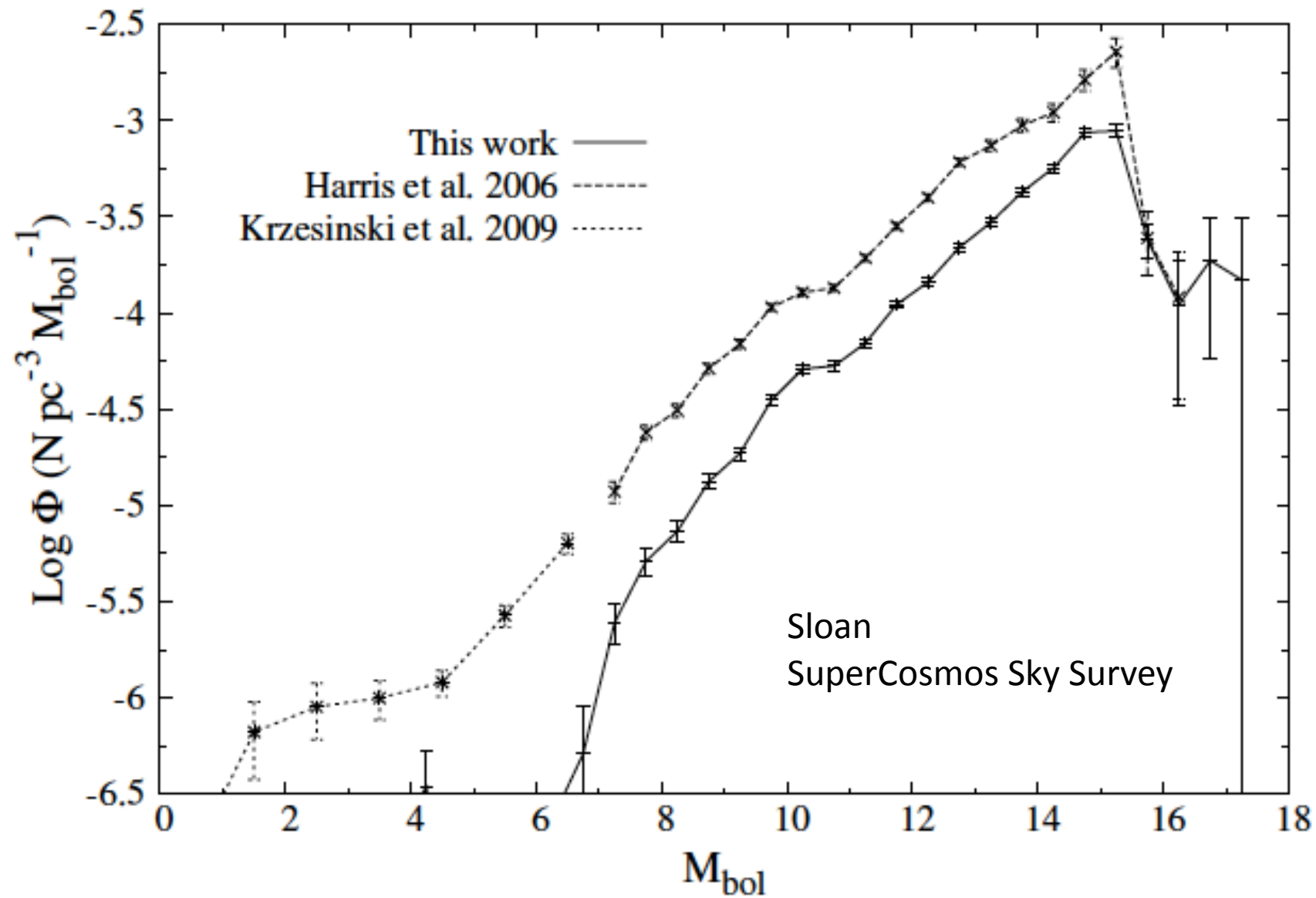
Surveys are more and more accurate and significant

Sample of WD:  
High precision LF



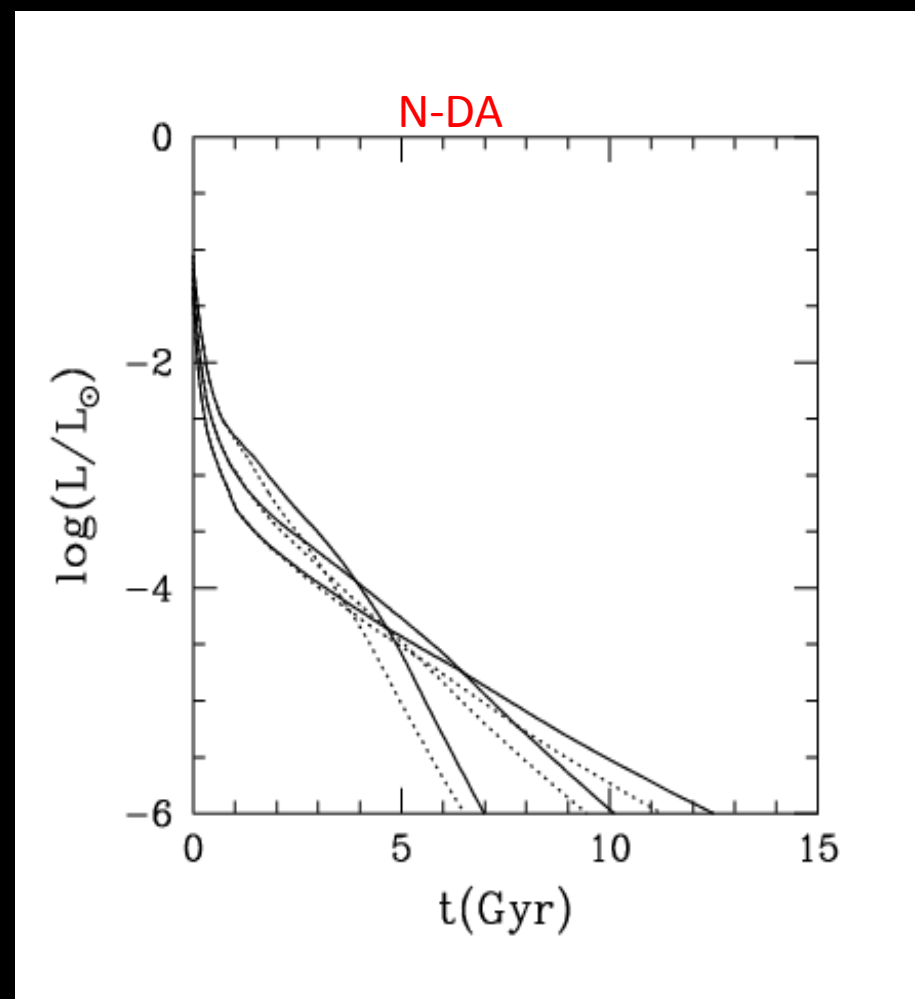
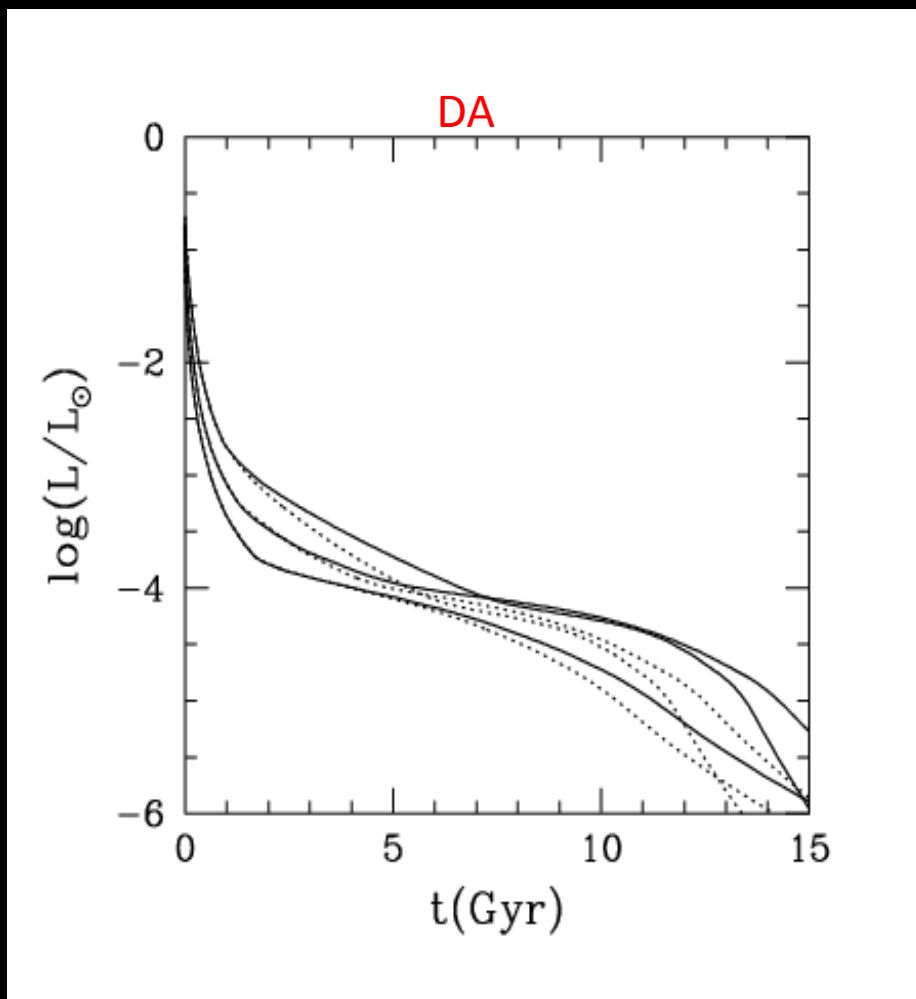
$$n(l) \propto \langle \tau_{\text{cool}} \rangle \int_{M_i}^{M_{\text{max}}} \Phi(M) \Psi(\tau) dM$$

Isern & Garcia-Berro'08



Rowell & Hambley'11

Luminosity versus time  
(dotted lines without sedimentation)



# DA, non-DA influence

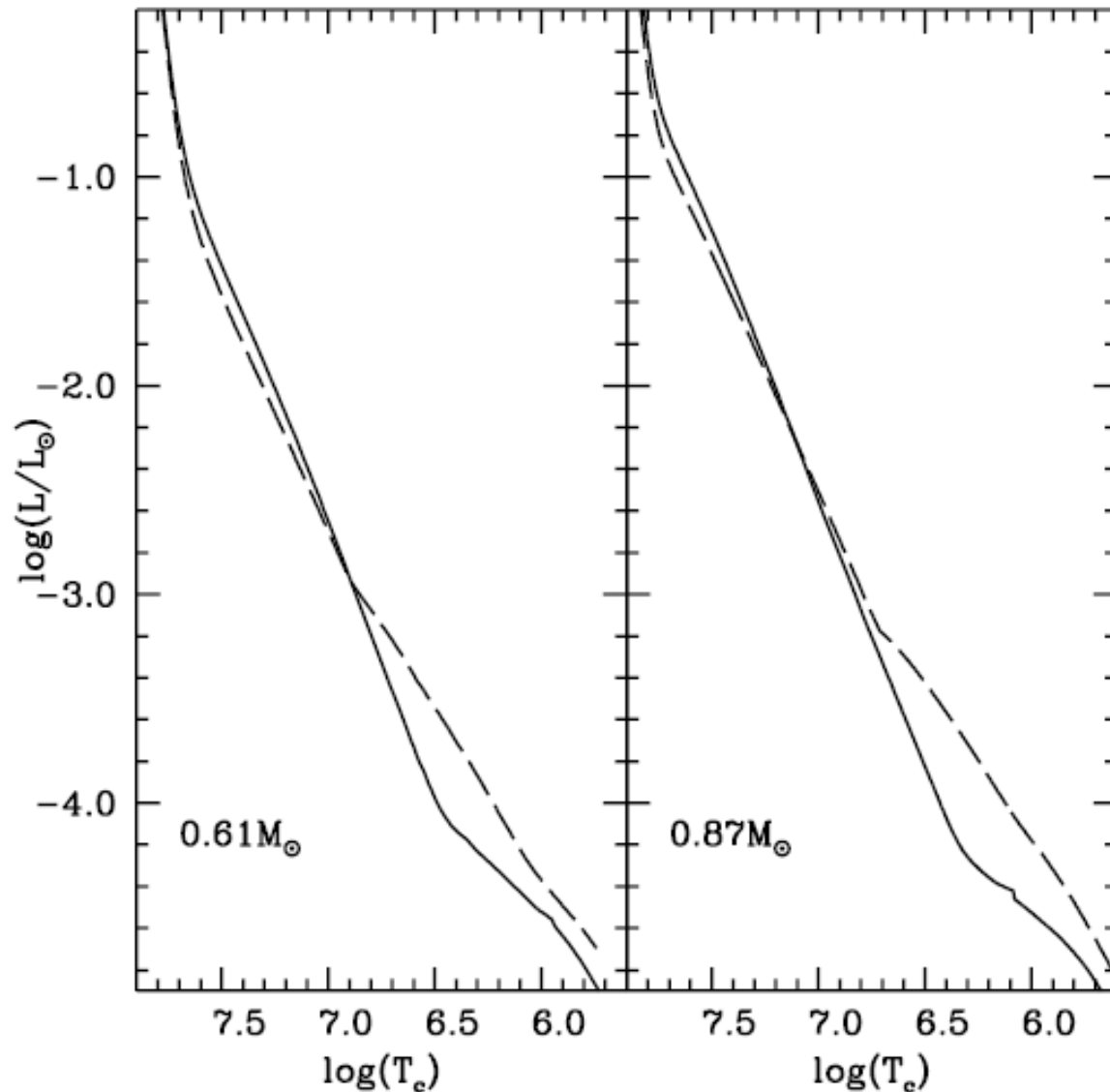


Fig. 1.—  $L - T_c$  relationships for our  $0.61$  and  $0.87 M_{\odot}$  WD models (with phase separation not included). Solid lines denote H-atmosphere models, dashed lines He-atmosphere ones.

Assume that:

$$L = g T_C^{\gamma}$$

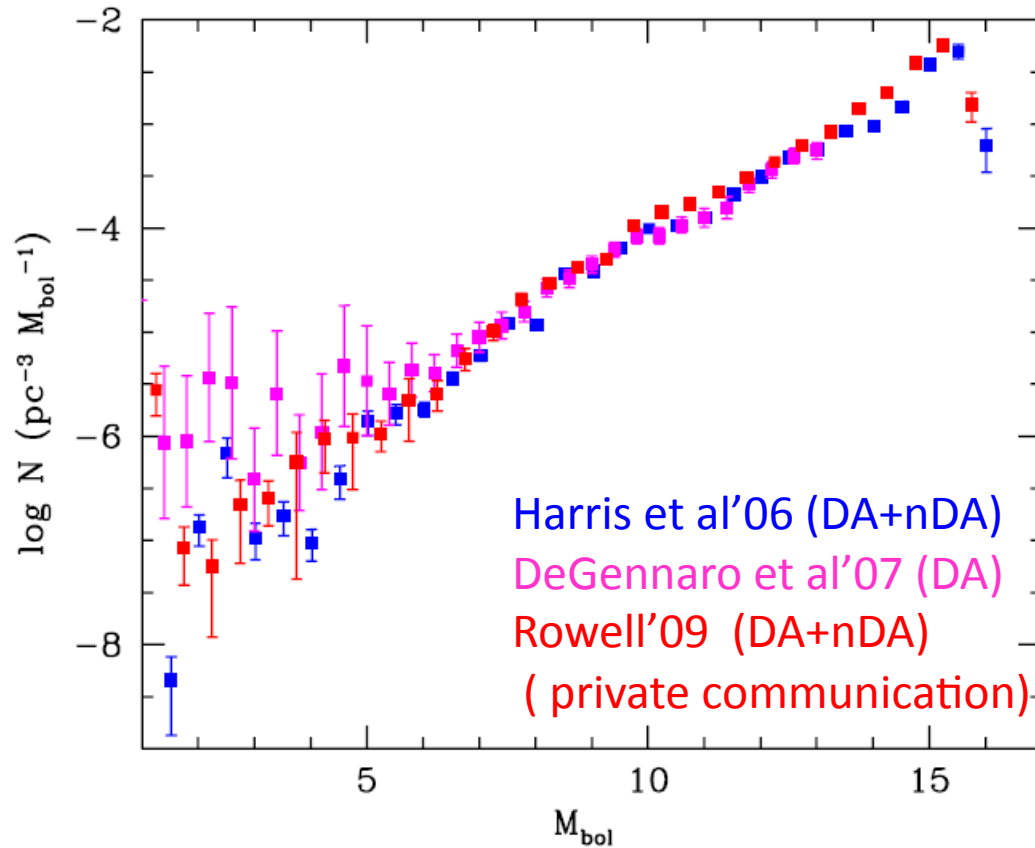
From the figure  
we see that:

$$\gamma_{DA} \approx \gamma_{nDA}$$

in the range

$$-3 \leq \log L \leq -1$$





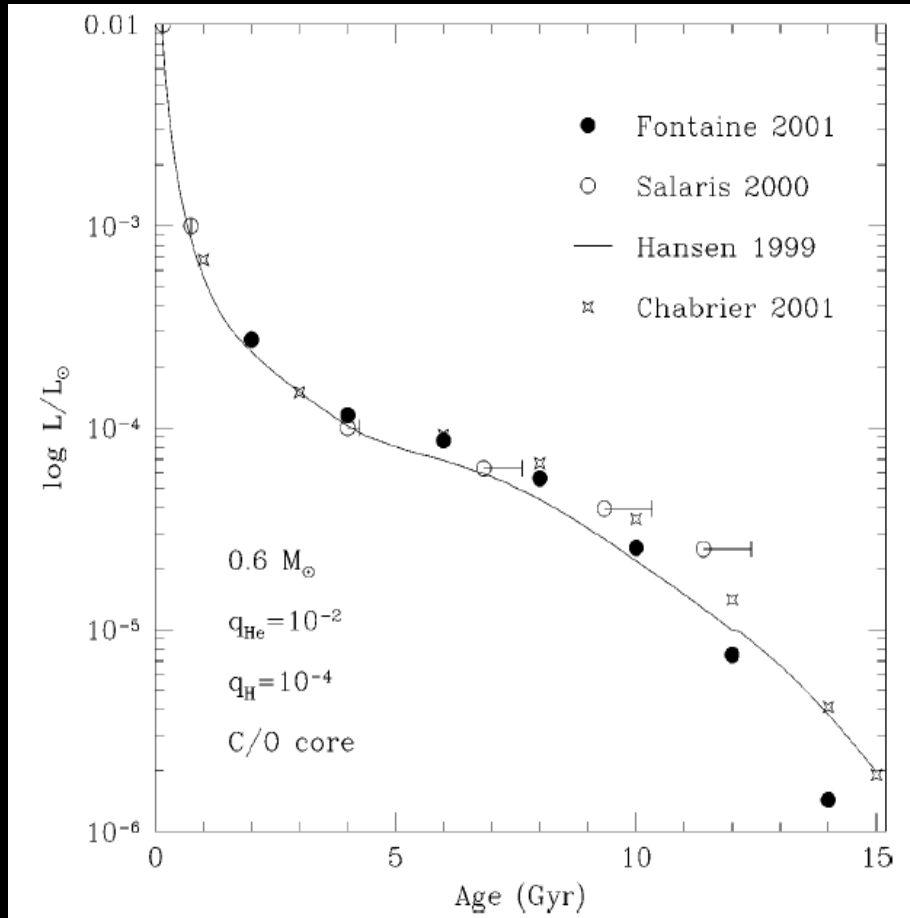
# Since  $\Upsilon_{DA} \approx \Upsilon_{nDA}$  the luminosity function of Das and nDAs coincide after normalization

$$L \approx -\frac{dU}{dt} \approx -C_V \frac{dT_C}{dt} \quad (\text{we neglect the compression term})$$

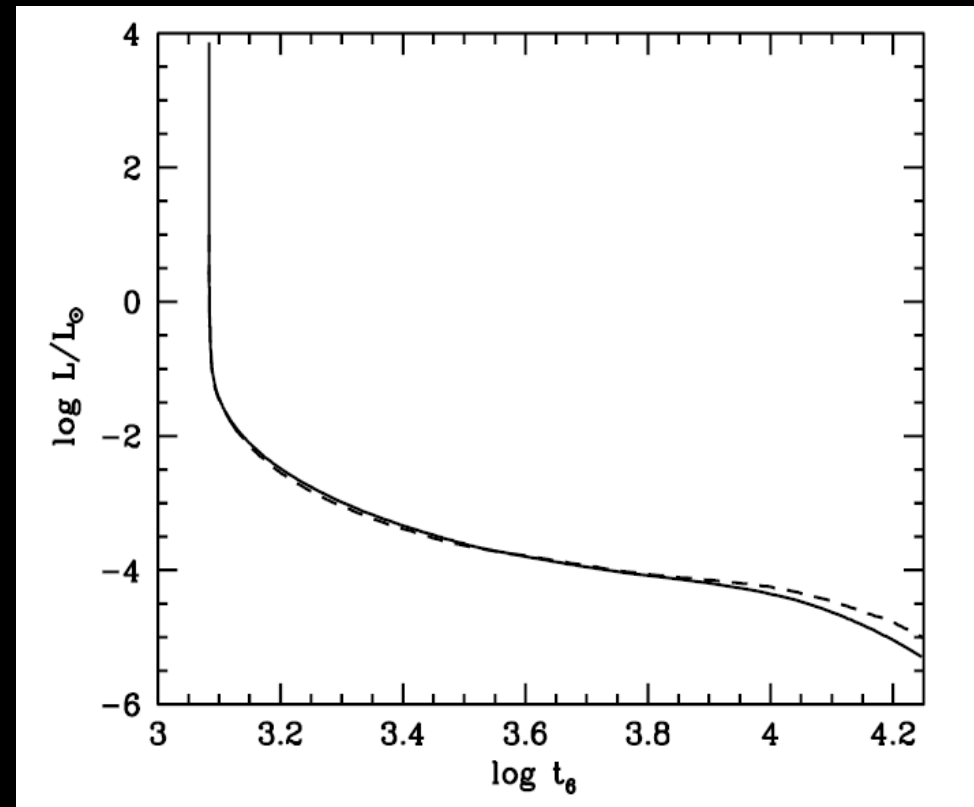
$$\frac{dL}{dt} = \gamma g T_C^{\gamma-1} \frac{dT_C}{dt} \quad (\text{from the L-}T_C \text{ relationship})$$

$$N_{WD} \propto \dot{L}^{-1} = -\frac{L}{dL/dt} = \frac{C_V}{\gamma g} T_C^{1-\gamma}$$

# Comparison between cooling models

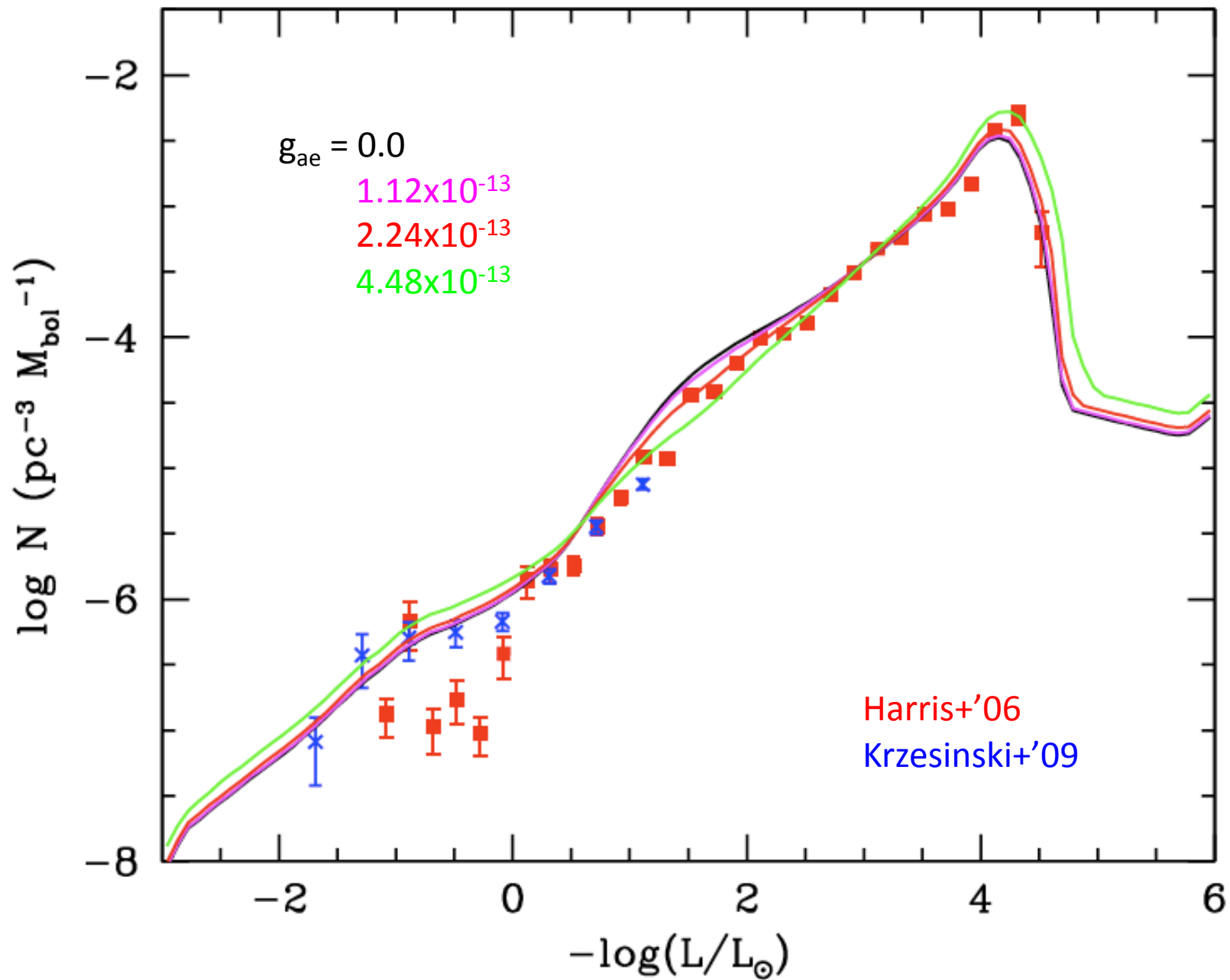


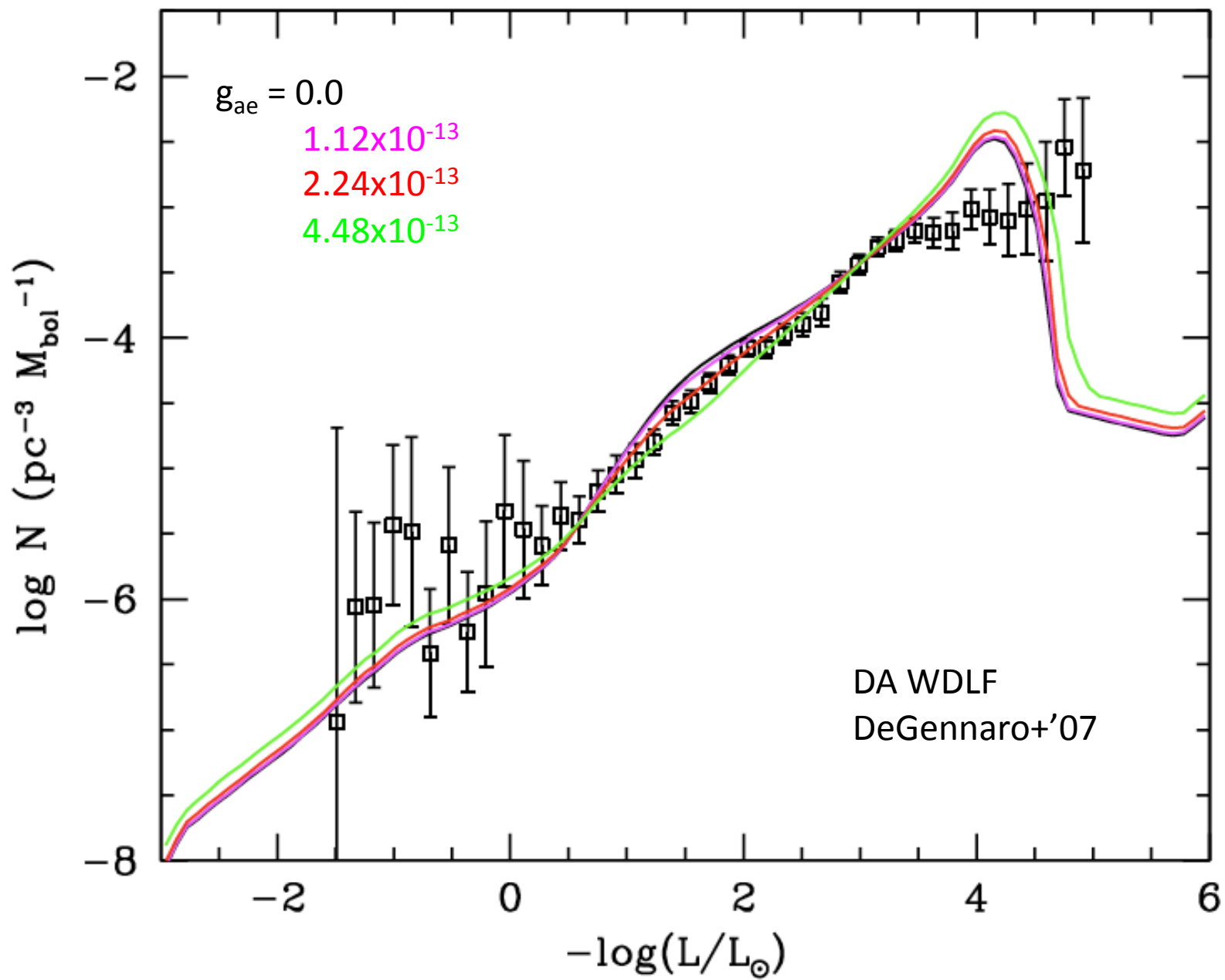
Hansen & Liebert'03



—: Renedo et al 2010

---- : Salaris et al 2010





# GAIA mission (2013-2018)



400,000 WD

# Large Synoptic Survey Telescope (LSST)

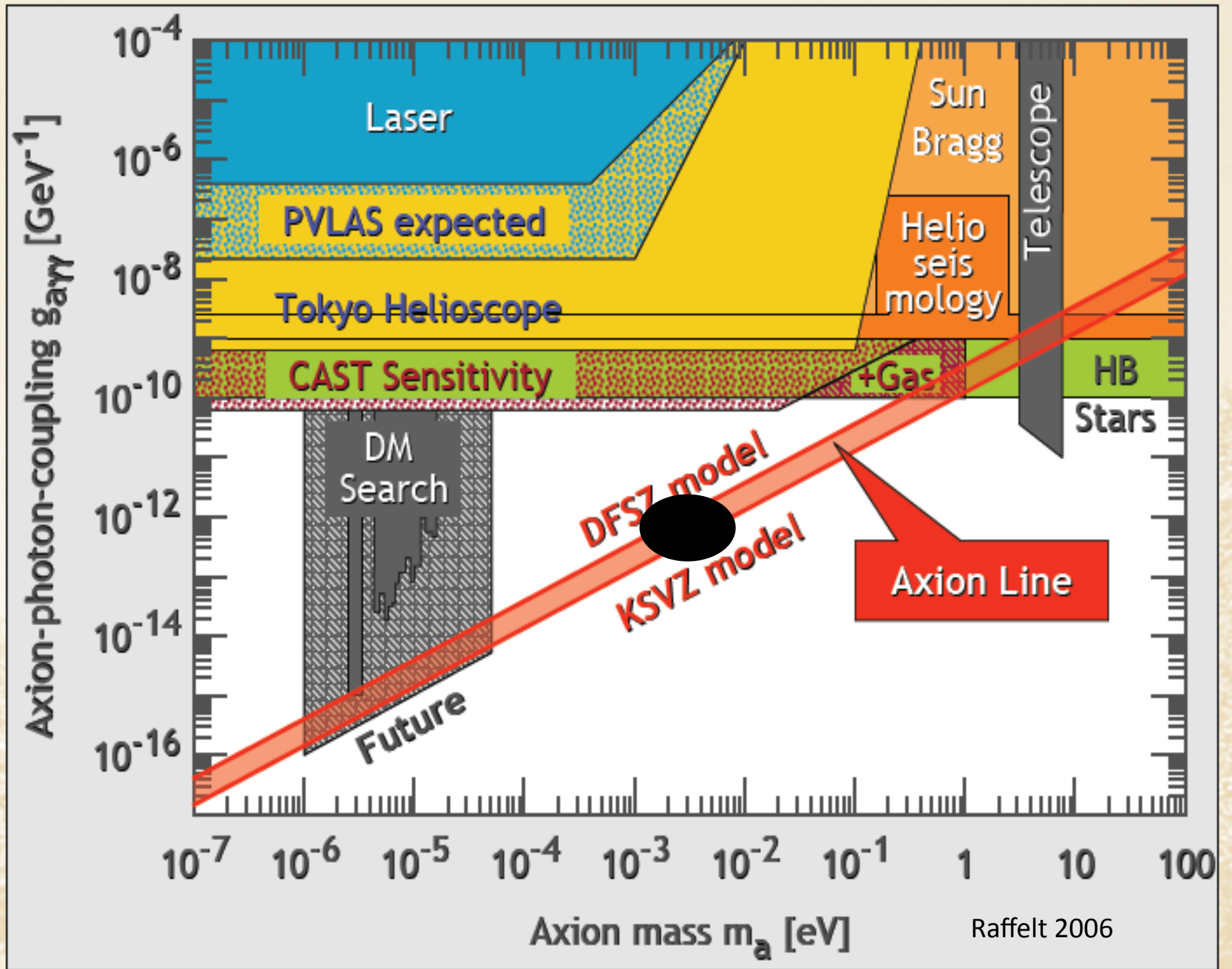


First light: 2015

Start Science: 2017

50,000,000 WD

$r > 27.5$  mag



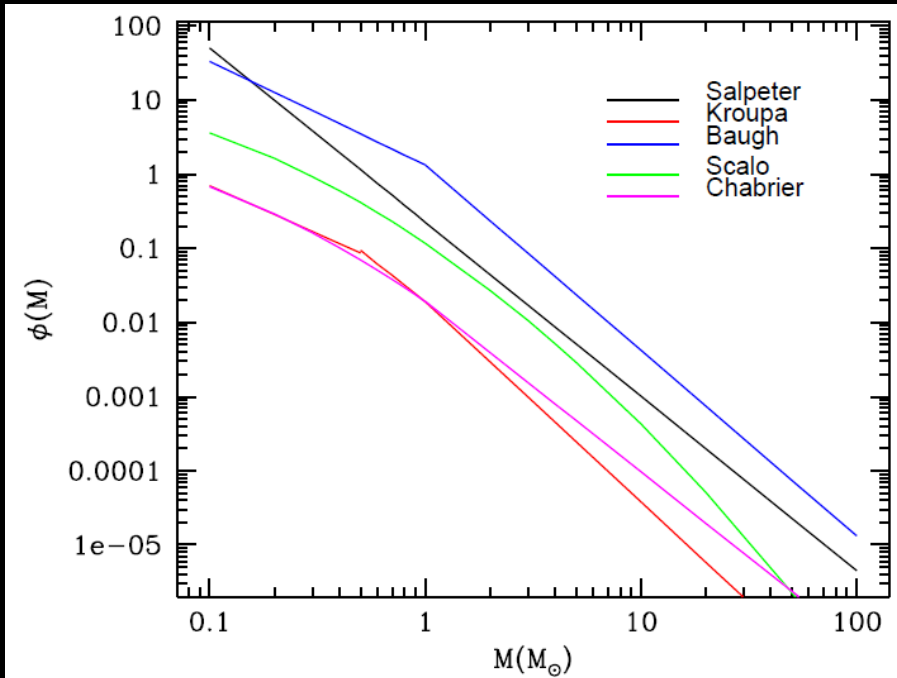
## Conclusions:

- # Because of their simplicity, WDs are excellent complementary laboratories for testing new physics.
- # The recent luminosity functions and the measurement of the secular drift of the pulsation period of DAV suggest that WDs cool down more quickly than expected .
- # Axions or light bosons able to couple to electrons could account for this discrepancy (  $g_{ae} \sim 2 \times 10^{-13}$  )
- # The results seem robust (for the moment) but more refinements are needed:
  - \* Extend the observational LF to high and low luminosities
  - \* Obtention of the LF for massive white dwarfs
  - \* Improvement of the cooling models. Envelope is crucial
  - \* Role of binaries
- # This method can be used in other problems

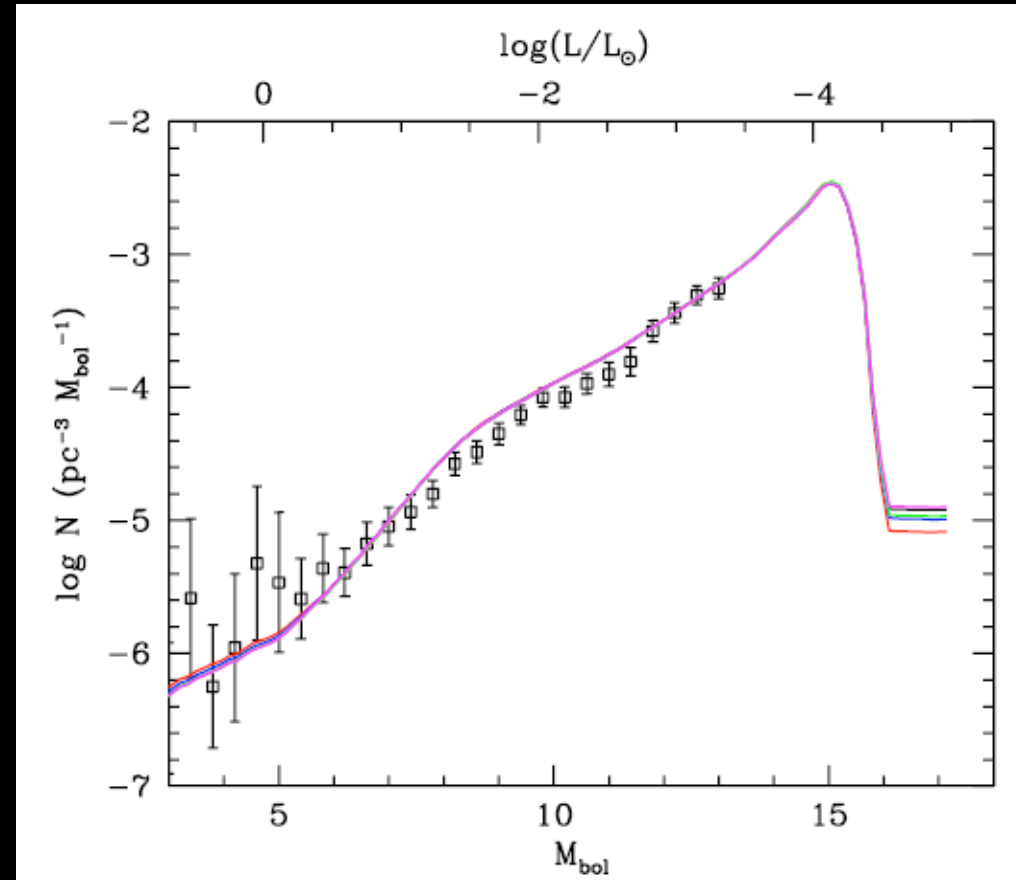
**GAIA & LSST can provide the necessary precision  
& accuracy**



# Dependence on the IMF

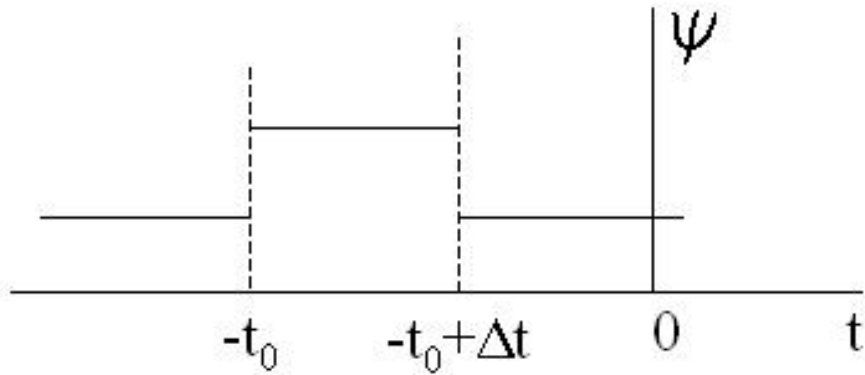


SFR=1 and the age=11 Gyr



The WDLF is not very dependent on the IMF as far as low mass stars are effectively produced.

# Influence of the SFR

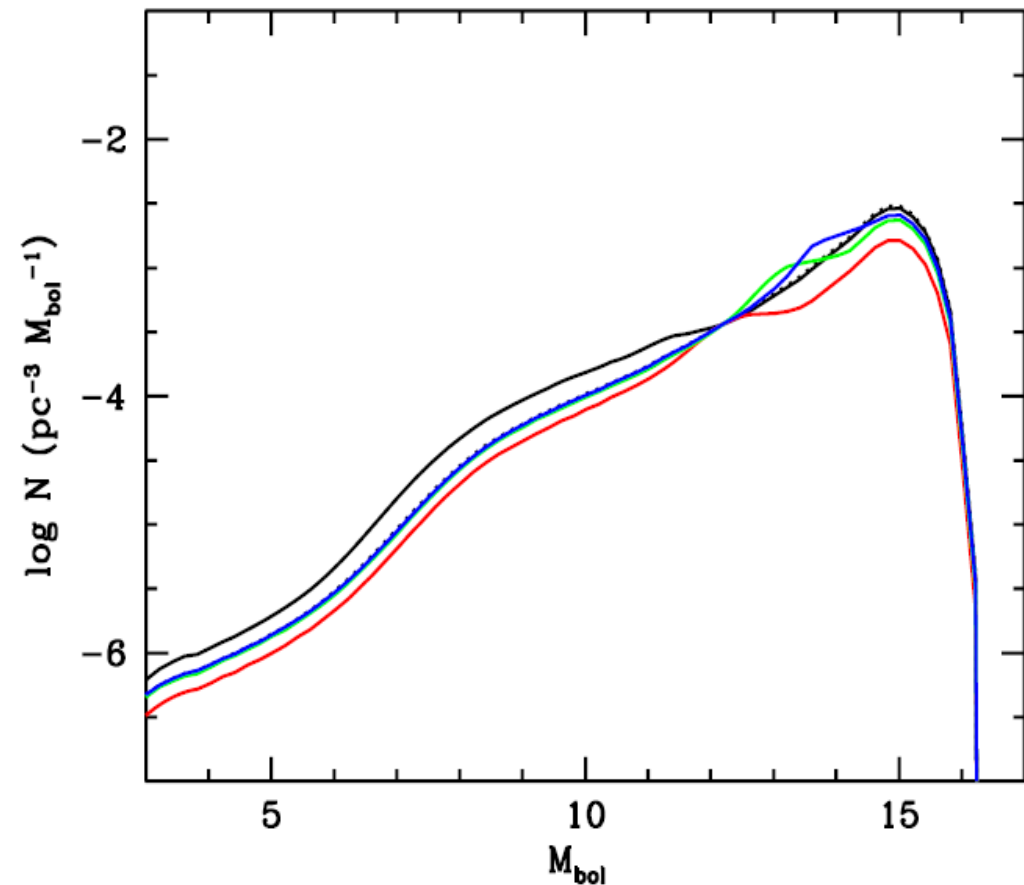


If the peak coincides with the normalization (red line) the bright branch falls below the standard

$T_0$	Color
0 (no bump)	Black dotted
-1	Black
-2	red
-3	Green
-4	Blue

$$\psi = 3, \text{ if } t_0 < t < t_0 + \Delta t$$

$$\psi = 1, \text{ if } t < t_0 ; t > t_0 + \Delta t$$



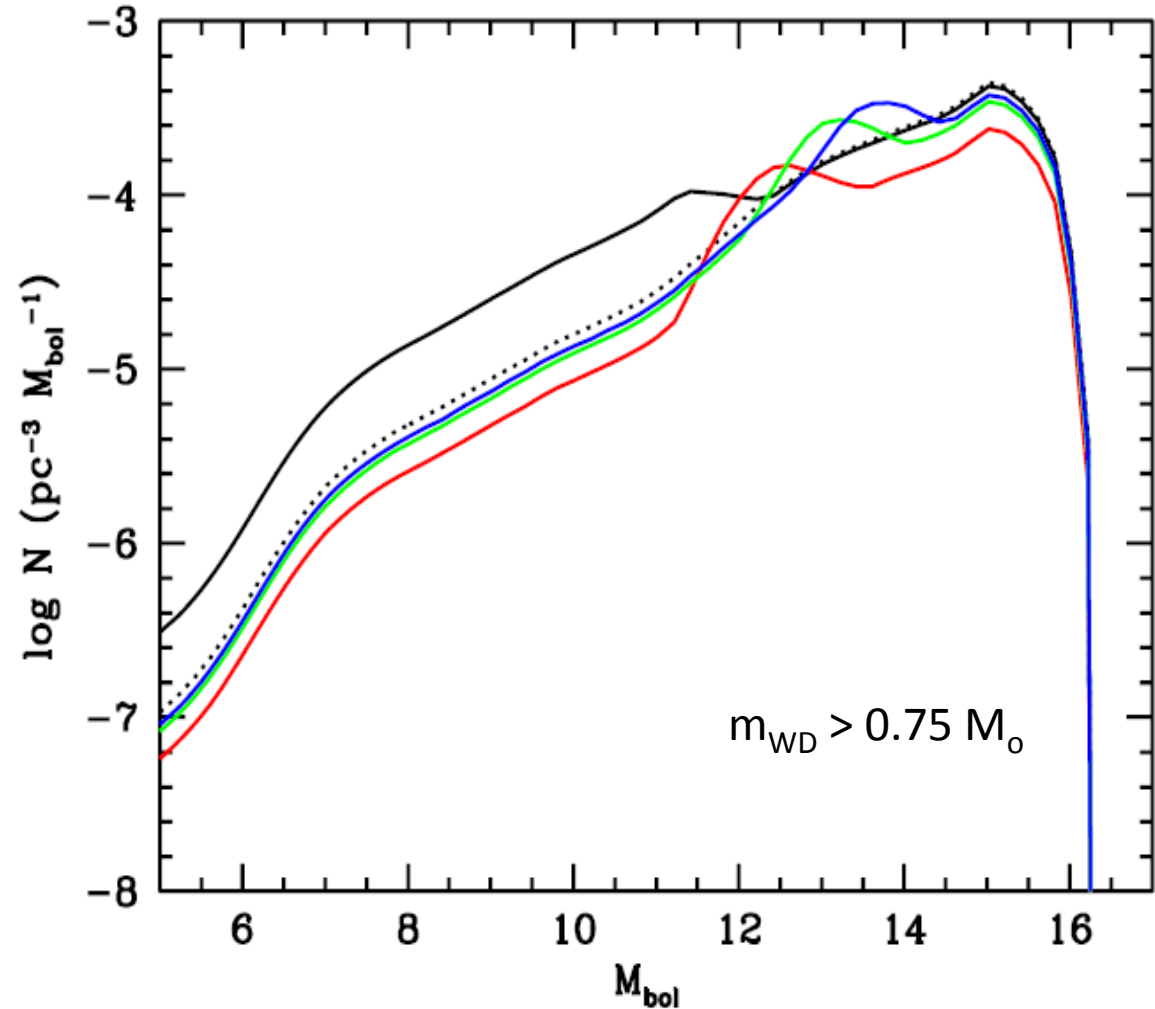
$$n(l) = \int_{M_{\min}}^{M_{\max}} \Phi(M) \Psi(T_{gal} - t_{cool} - t_{SP}) \tau_{cool} dM$$

In the case of massive WD

$$t_{SP} \ll t_{cool}$$

$$n(l) \propto \Psi(T_{gal} - t_{cool})$$

The luminosity function of massive WD closely follows  
The SFR  
Irregularities are detectable!



# Influence of binaries:

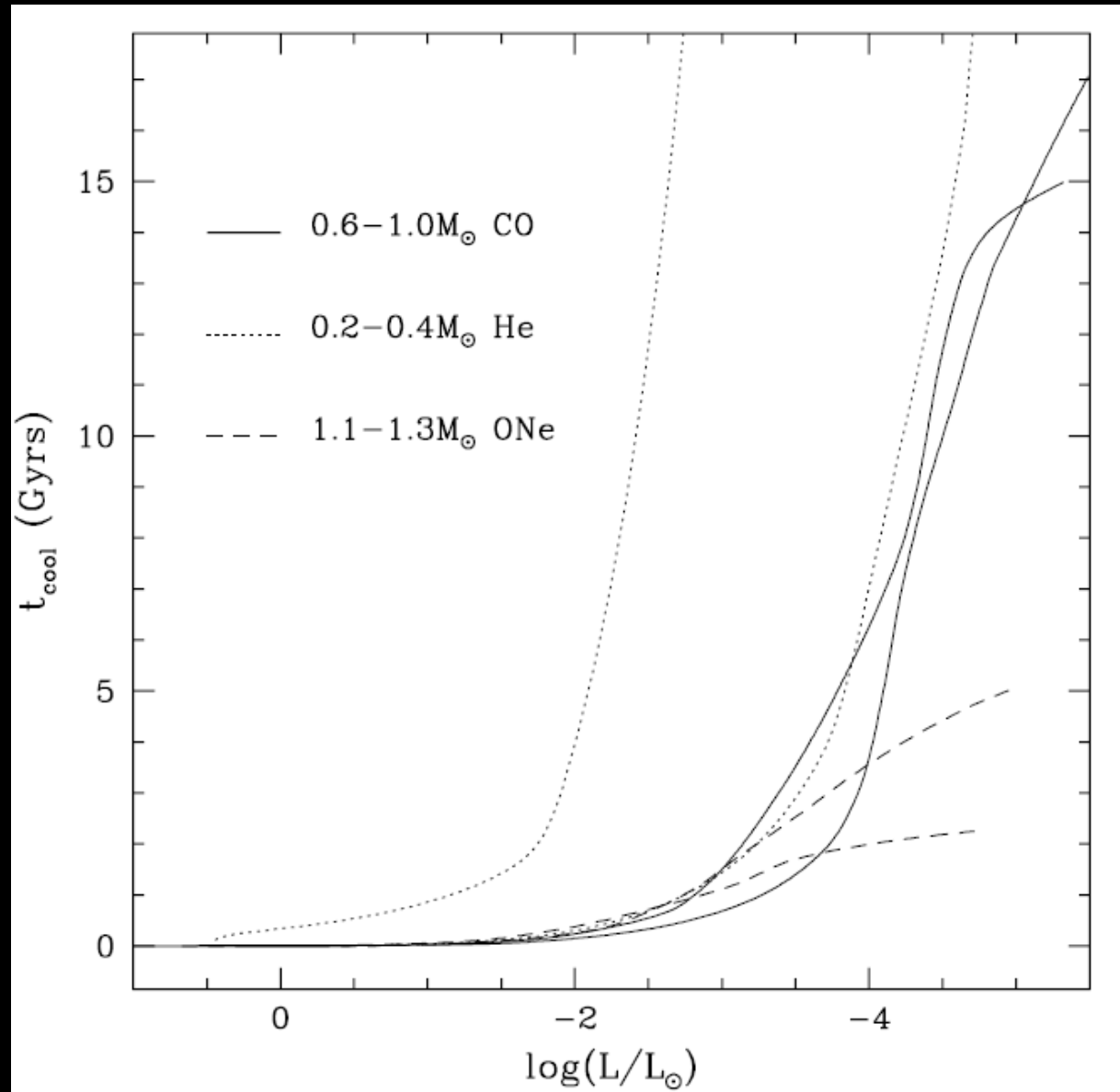
# Presence of He-white dwarfs

# Mergers

# Tidal heating

# Non resolved binaries

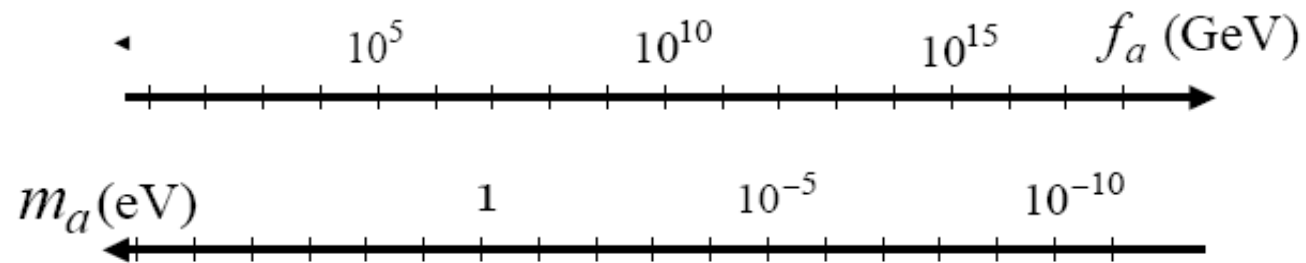
Contamination by He-WD



# The axion case

- Axions were proposed as a solution to the strong CP problem
  - KVSZ model -> Axions couple to hadrons & photons
  - DFSZ model -> Axions also couple to electrons
- Coupling is determined by the Peccei-Quinn scale  $f_a$  which is related to the mass of the axion:  $m_a = 6.0 \text{ eV} \cdot (10^6 \text{ GeV}/f_a)$
- Experiments have failed to detect axions
- Constraints from astrophysical arguments
  - Solar properties
  - Red giants (HB & AGB stars)
  - Core collapse supernovae
  - Cosmological considerations

## The remaining axion window



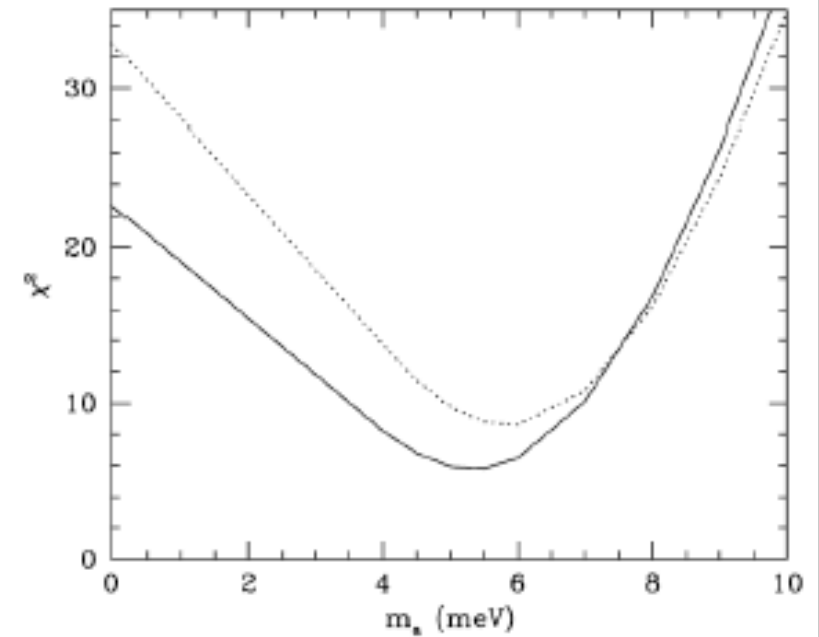
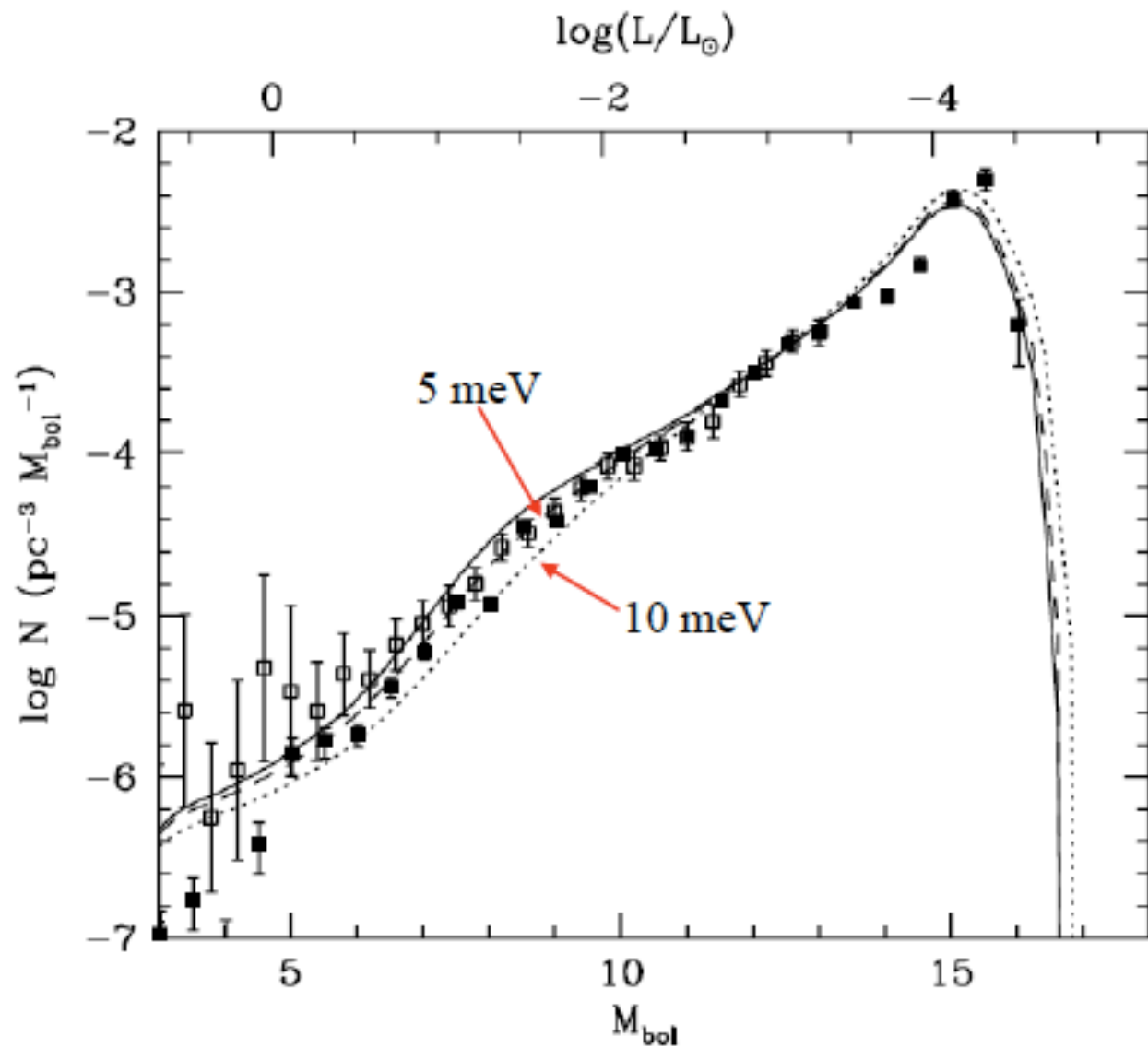
laboratory  
searches

stellar  
evolution

cosmology

$10^{-2} - 10^{-6}$  eV

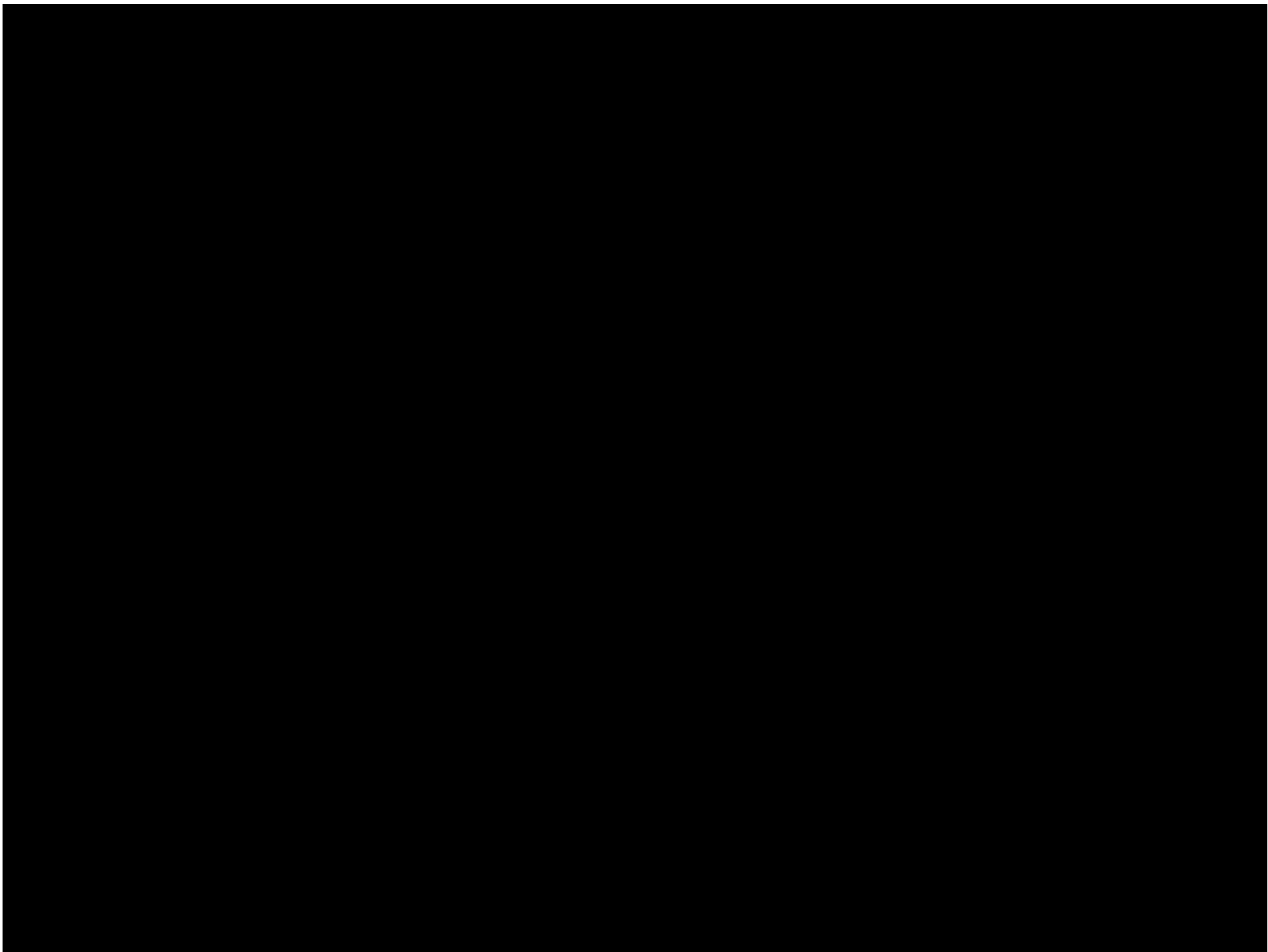
For these masses, axions can freely escape from stars  
They can be treated as a sink of energy



**Isern et al'08**

The best fit is obtained for  $m_a \cos^2 \beta \sim 5$  meV

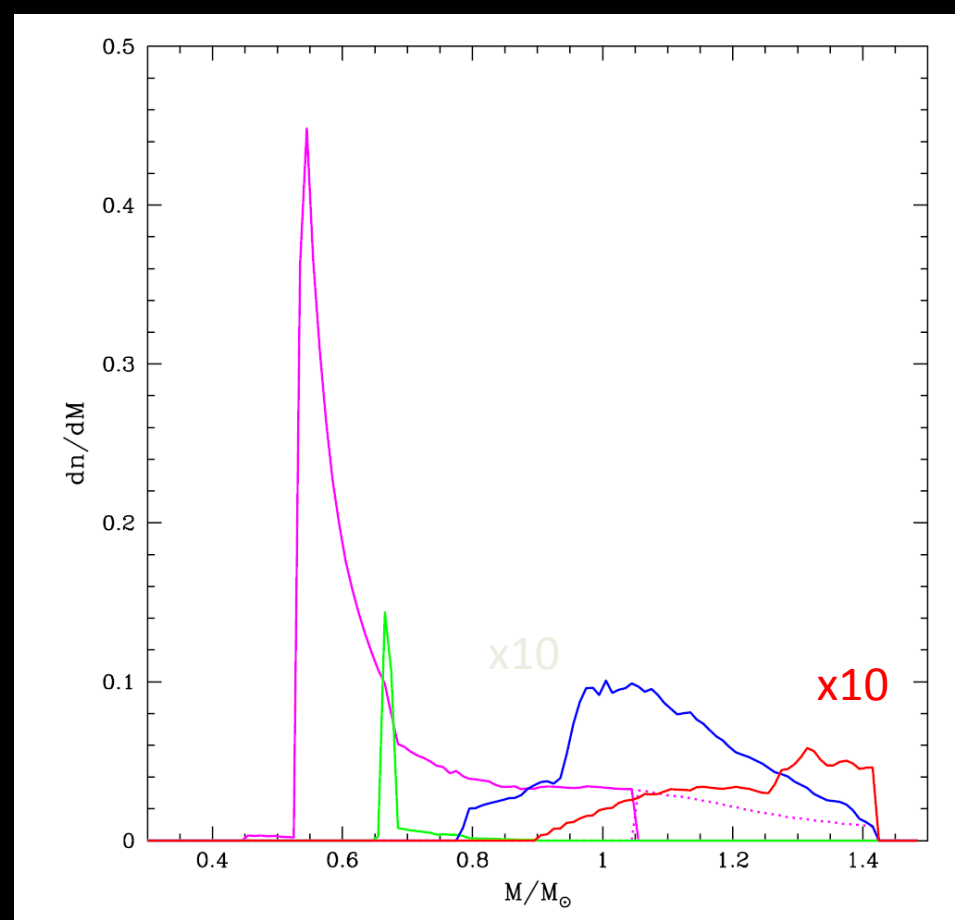
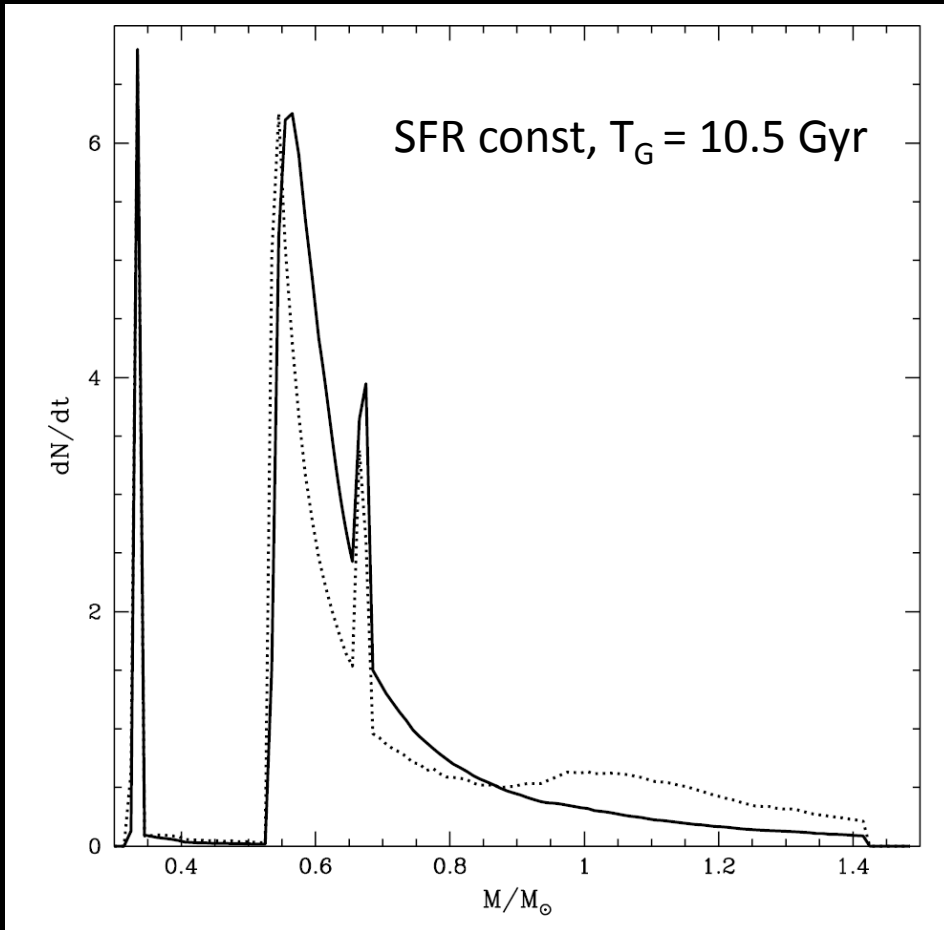




# Birthrate calculation

Isern et al,  
Thermonuclear Supernovae,  
Ed. Ruiz-Lapuente, Canal, Isern,  
Kluwer p. 127 (1997)

- Only evolutionary channels in which RLOF occurs when the envelope is convective
- Models obtained with FRANEC. Solar metallicity
- WD cooling models from Salaris et al 2000
- Catalán et al (2008) IFMR
- Common envelope treatment: Iben & Tutukov (1984)
- Magnetic breaking
- Salpeter's IMF for the primary,
- $F(q) \propto q$ ;  $q = M_2/M_1$
- Distribution of initial separations:  $H(A_0) \propto 1/A_0$
- During the merging ALL the mass of the secondary is transferred to the primary

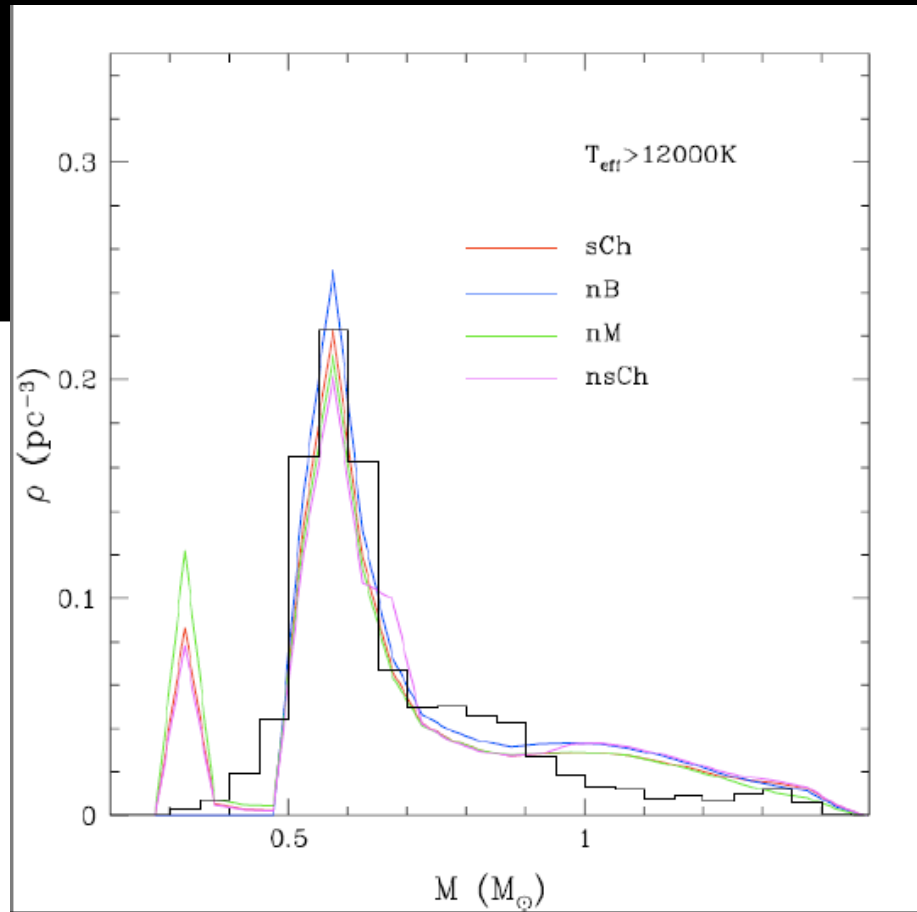
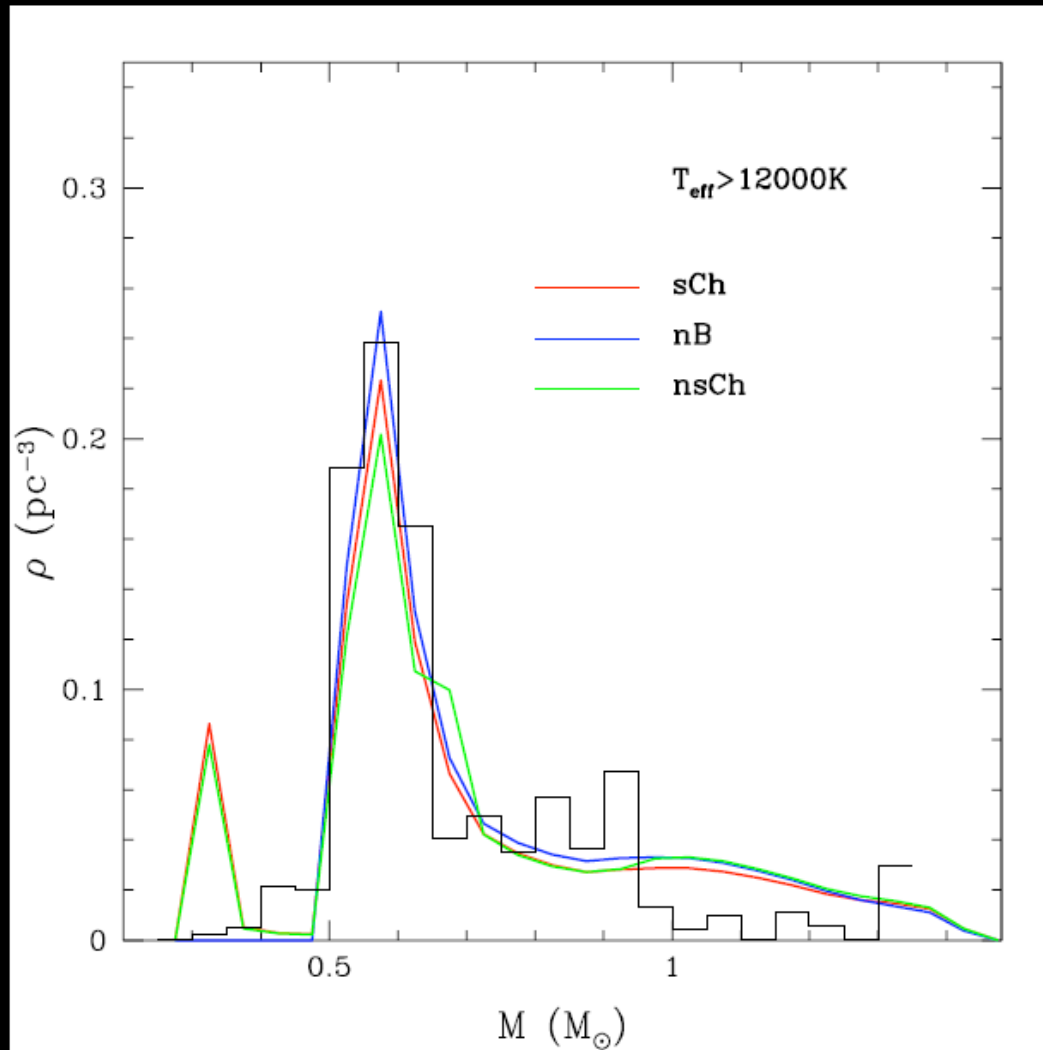


Single:  
M=1.05

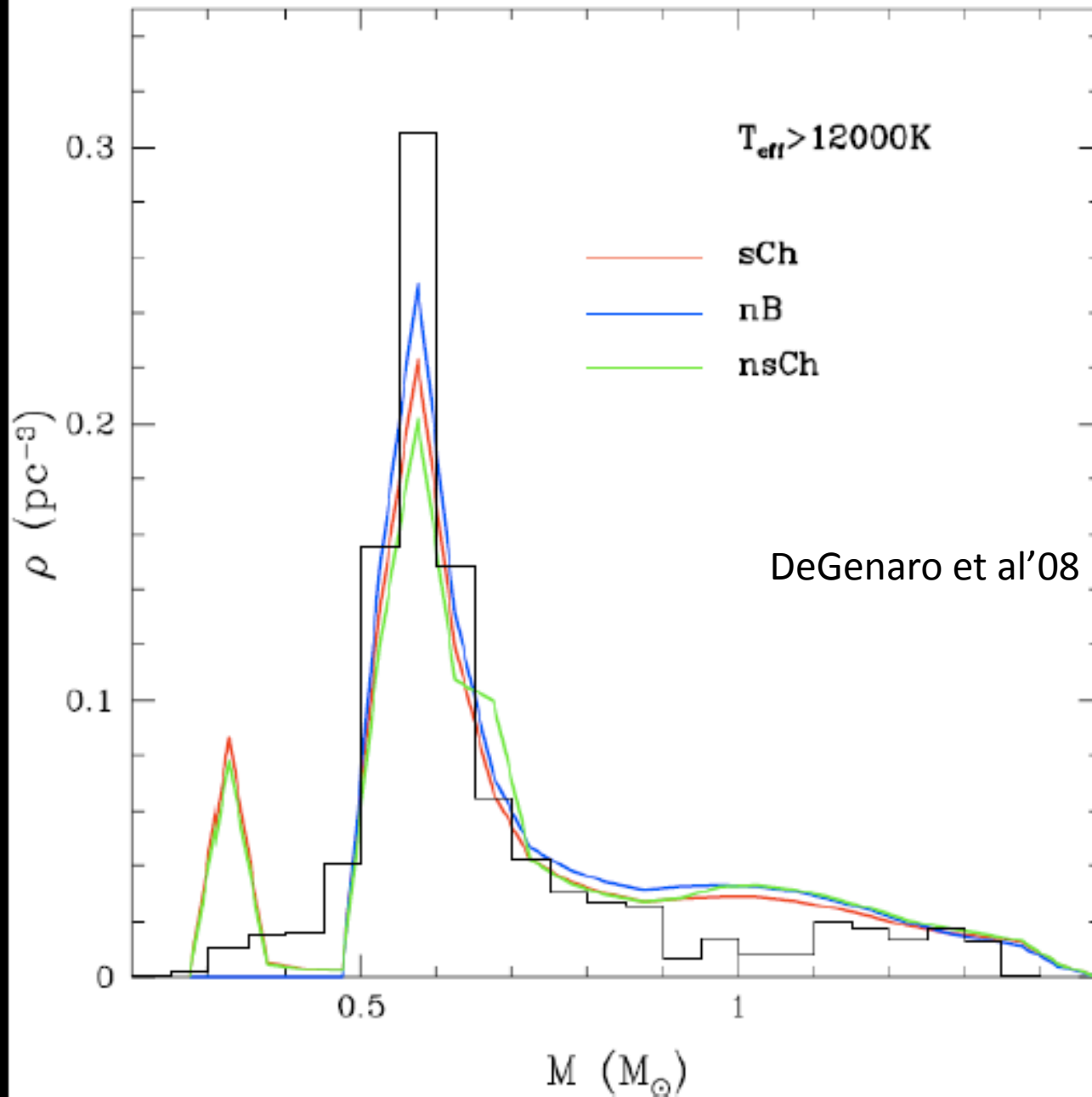
He + He  
M=0.65  
 $n_{\text{HeHe}} = 58\%$   
 $n_{\text{CO}} = 42\%$

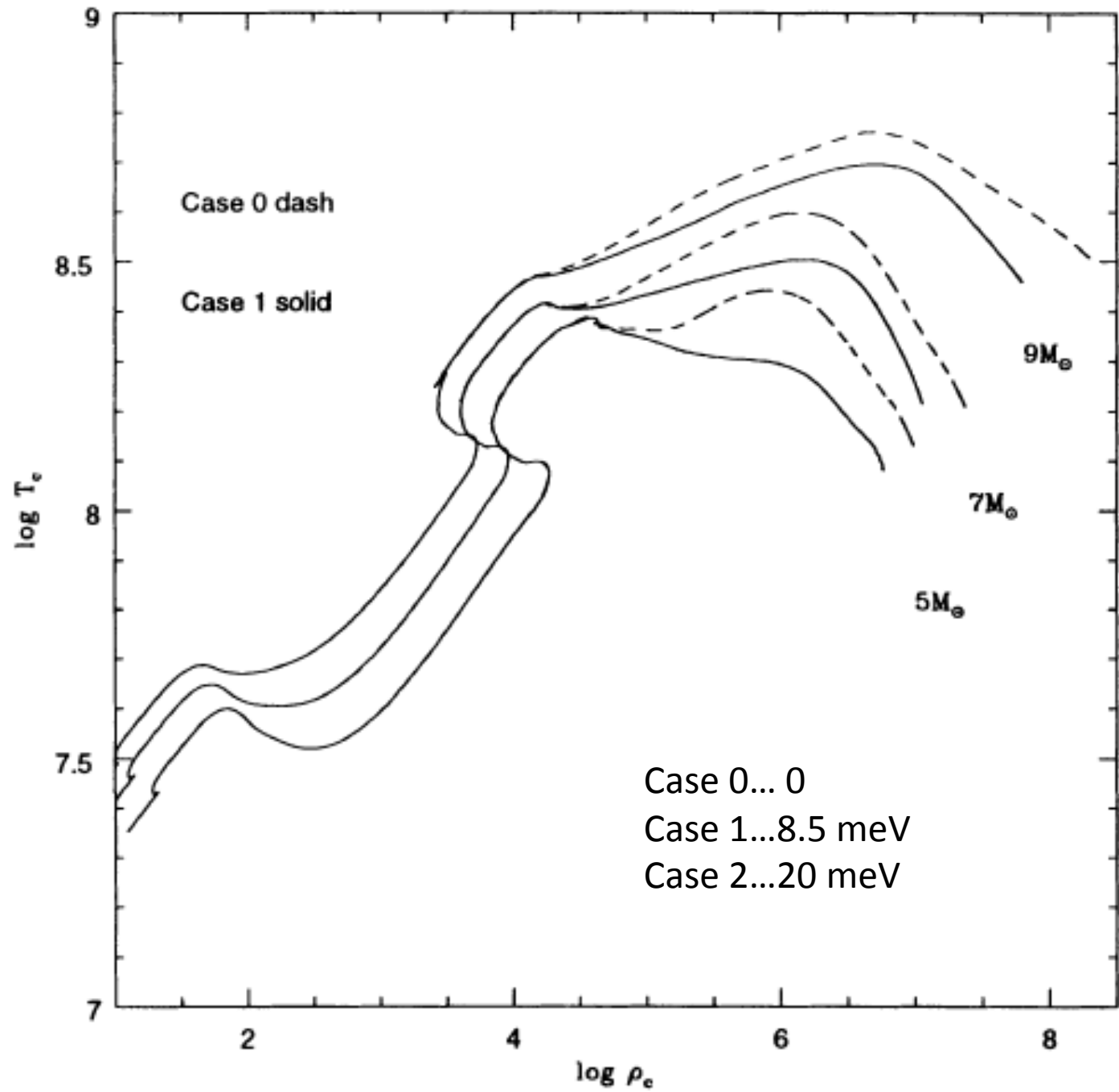
CO+He: ~ 20%  
CO+Co: ~ 15%

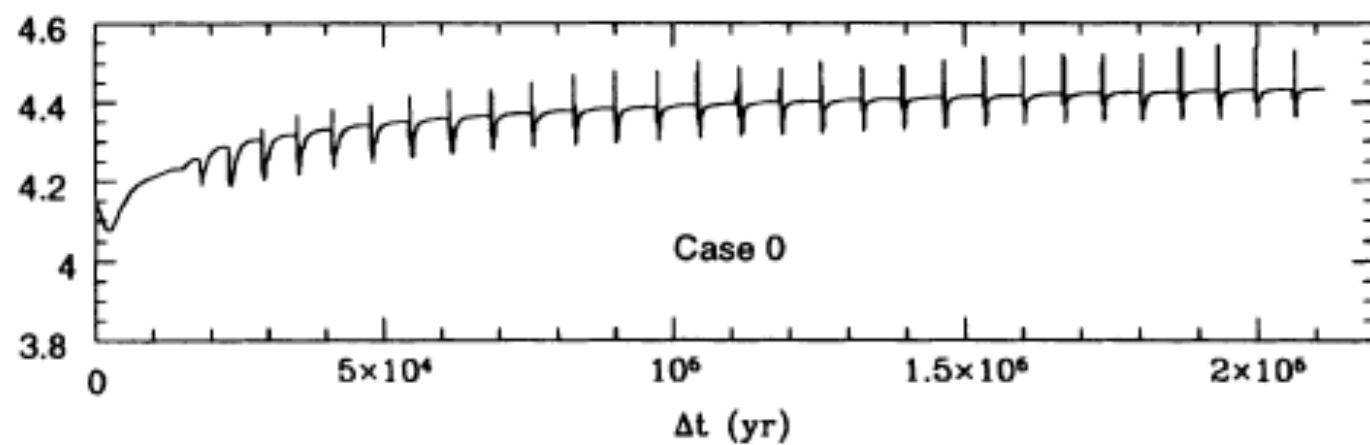
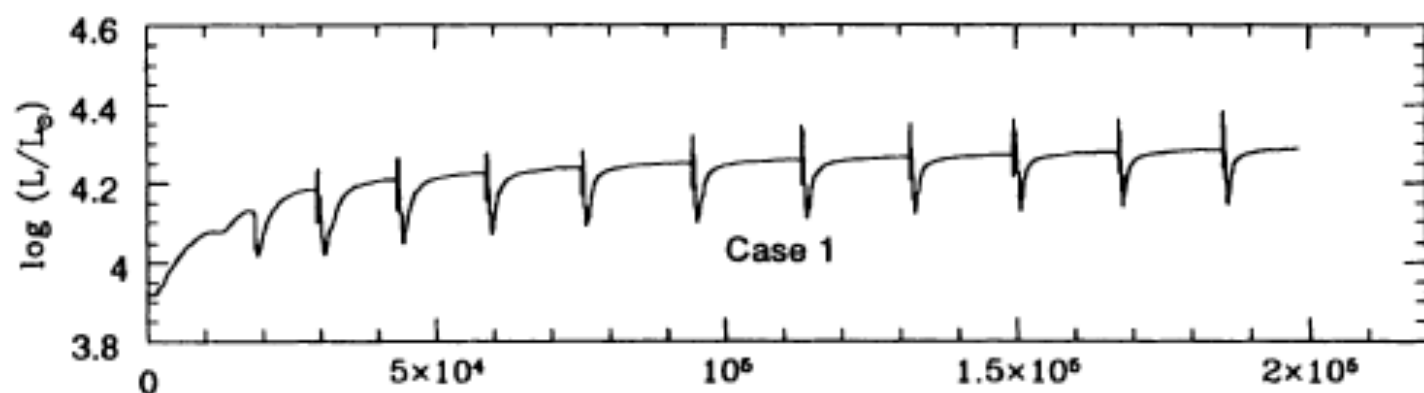
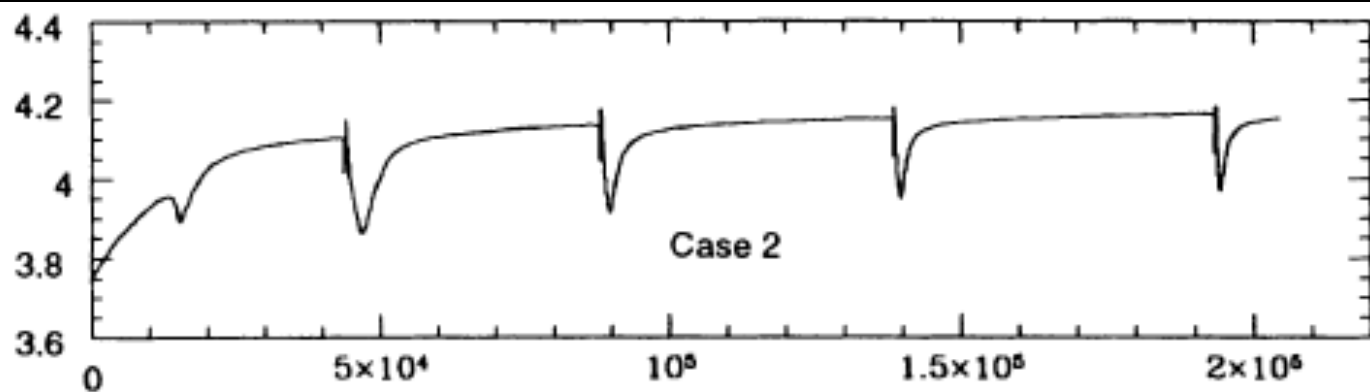
Liebert et al'05  
Bin average 0.05

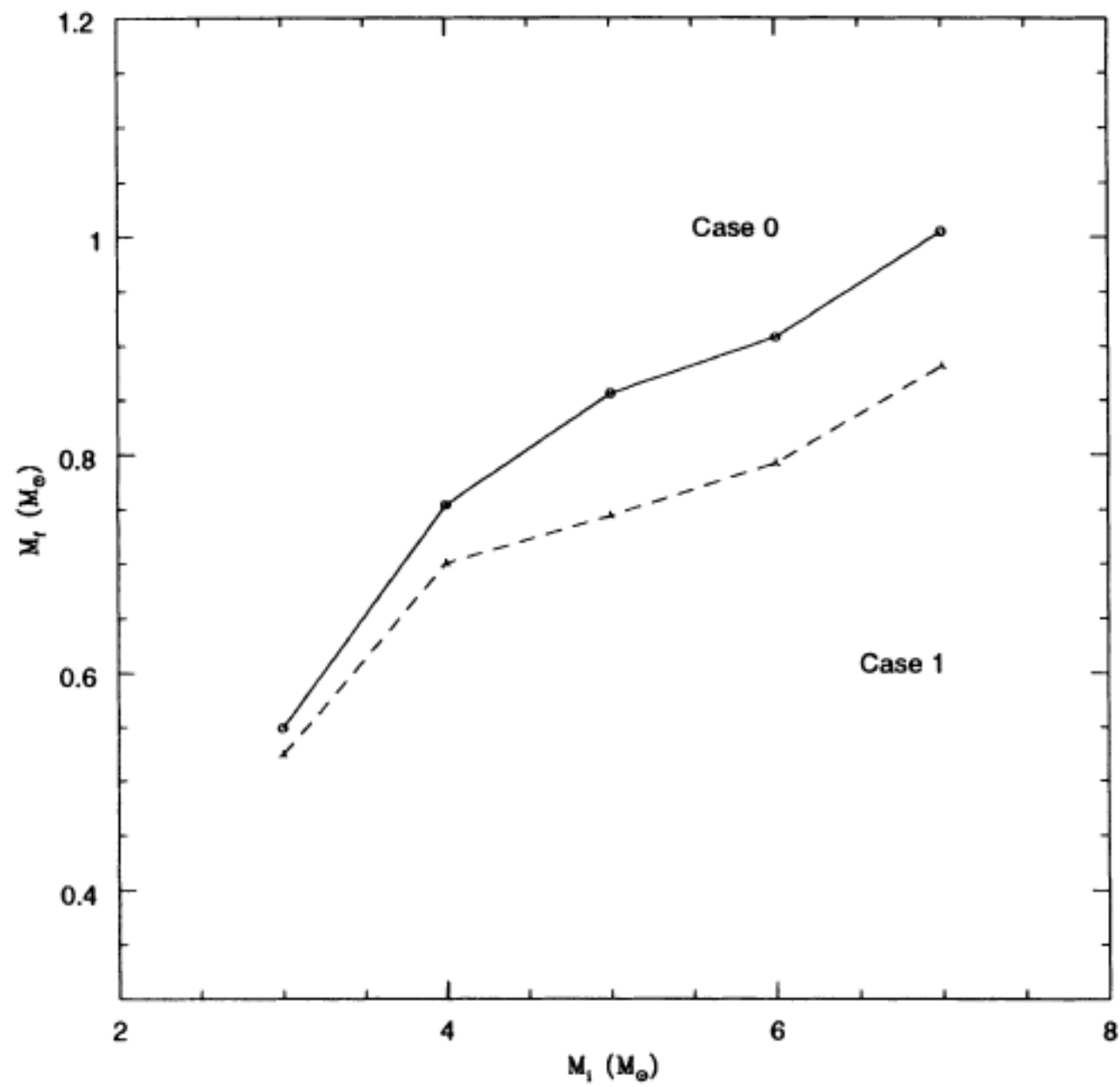


Gaussian errors included



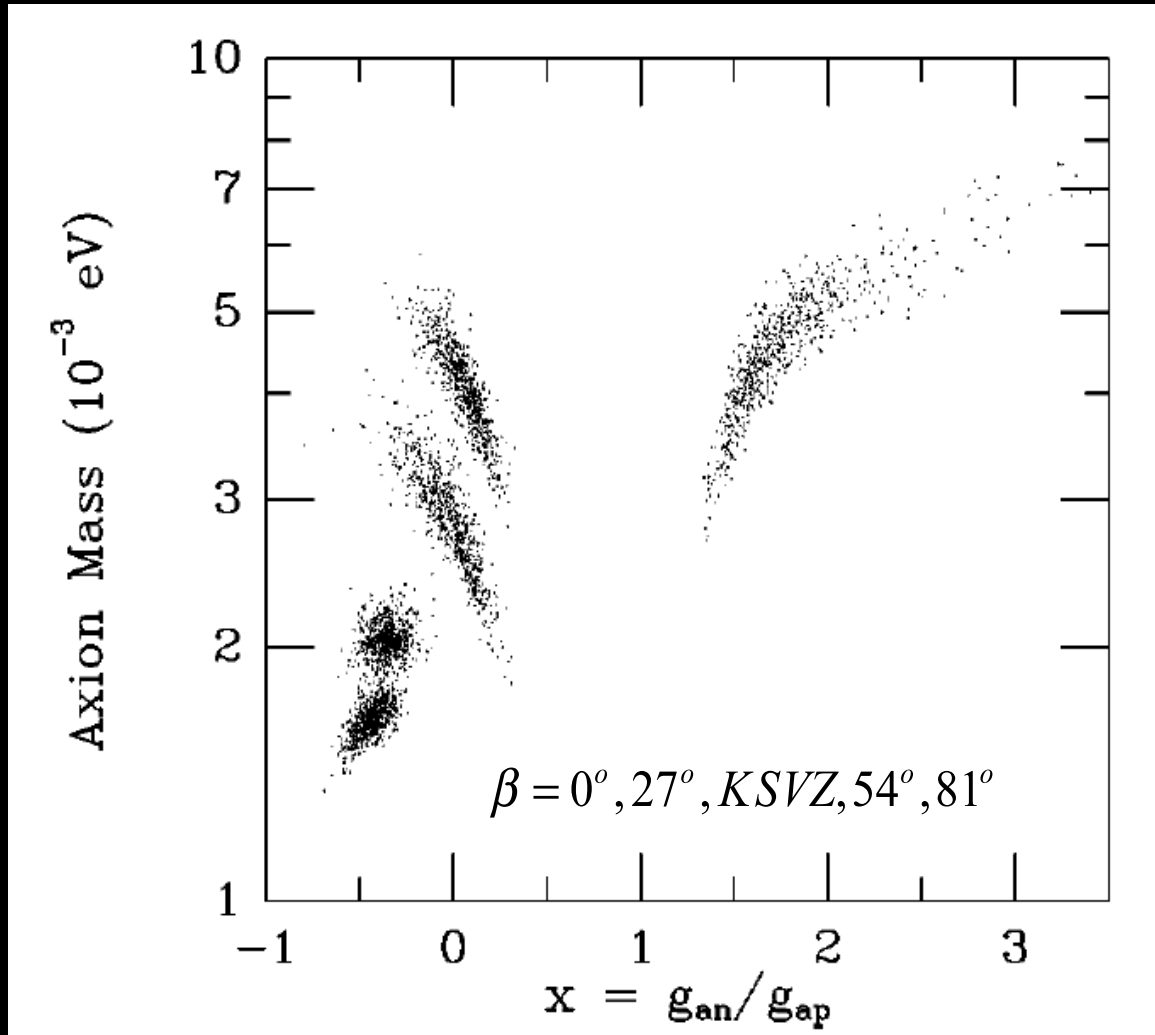








# Influence on core collapse supernovae



Keil et al '97

Nucleon bremsstrahlung is dominant

Raffelt'06

$m_a(KSVZ) < 16$  meV

$m_a(DFSZ) ?$